

Department of Physics
Vivekanand College, Kolhapur (Autonomous)

Notice for Internal Examination in Physics for B.Sc. III


It is hereby informed that; students of B.Sc. III should note that their Internal Examination in Physics will be conducted as per following time table.

Date	Time	Class	Subject
Tuesday, 22/11/2022	02.30 to 03.30 PM	B.Sc. III	Paper V section I (Classical Mechanics)
Wednesday, 23/11/2022	02.30 to 03.30 PM		Paper V section II (Quantum Mechanics)
Thursday, 24/11/2022	02.30 to 03.30 PM		Paper VI section I (Nuclear and Particle Physics)
Friday, 25/11/2022	02.30 to 03.30 PM		Paper VI section II (Mathematical Physics)

Nature of Question Paper

- Q.1) Select correct alternative (5 Marks)
Q.2) Long answer type question (10 Marks, Attempt any One)
Q.3) Short answer type question (5 Marks, Attempt any One)
Total Marks: 20 Marks




HOD, Physics
Head of the
Department of Physics
Vivekanand College, Kolhapur

Shri Swami Vivekanand Shikshan Sanstha's

Vivekanand College, Kolhapur

(Autonomous)

Department of Physics

Internal exam (2022-23)

B.Sc.III Sem V

Attendance Sheet

Roll No.	Name Of The Student	Signature			
		22/11/2022	23/11/2022	24/11/2022	25/11/2022
8201	Bhingardeve Dhiraj Prakash	Dhiraj	Dhiraj	Dhiraj	Dhiraj
8202	Dongare Prathamesh Abaji	PADongare	PAADongare	PDongare	PAADongare
8203	Dongare Suyash Sanjay	Songare	Songare	Songare	Songare
8204	Gaikwad Rajnandini Ganesh	ggaiwad	ggaiwad	ggaiwad	ggaiwad
8205	Jadhav Sae Sandeep	Sai	Sai	Sai	Sai
8206	Jamadar Mahek Shakilahmed	Mahek	Mahek	Mahek	Mahek
8207	Kalkutki Shubham Babasaheb	JKalkutki	JKalkutki	JKalkutki	JKalkutki
8208	Kamble Anjali Bhagwan	Anjali	Anjali	Anjali	Anjali
8209	Kothawale Tejas Vikas	Alonee	Alonee	Alonee	Alonee
8210	Maner Aman Imtiyaj	Amance	Amance	Amance	Amance
8211	Padmakar Alok Narayan	ATejas	ATejas	ATejas	ATejas
8212	Patil Aaryan Pramod	aeeyan	aeeyan	aeeyan	aeeyan
8213	Shinde Vivek Janardan	Shinde	Shinde	Shinde	Shinde
8214	Shingade Aishwarya Deepak	Shingh	Shingh	Shingh	Shingh
8215	Singh Sapana Raviranjan	Sanske	Sanske	Sanske	Sanske
8216	Warke Shriyash Keraba	Warke	Warke	Warke	Warke
8217	Yadav Vedaja Ajay	Y.A.Yadav	Y.A.Yadav	Y.A.Yadav	Y.A.Yadav

Internal Examiner.....



Shri Swami Vivekanand Shikshan Sanstha's
Vivekanand College, Kolhapur (Autonomous)

Internal Examination 2022-23

PHYSICS-DSC -1001E

B.Sc. – III, Sem – V Classical Mechanics

Time: 30 Minutes

Marks: 20

Q. 1. LONG Answer Questions (Any one)

(10)

- 1) Derive Lagrange's equation from D'Alembert's principle.

OR

- 1) What do you mean by constraints. Explain its types.

Q. 2. SHORT Answer Questions

(10)

- 1) Write a note on Atwood's Machine.
- 2) Write a note on degrees of freedom.



Name :- Dhanashree Ananda Raval.
Class :- Bsc - III Div :- C Roll No :- 8208
Sem :- I.

INTERNAL EXAM

VIVEKANDA COLLEGE, KOLHAPUR.

75

Q.1. Answer the following questions.

1. Explain analogy of rotational motion with translational motion.

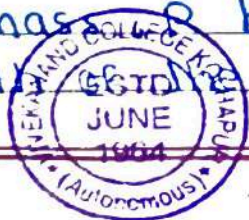
-
- i) It is noticed that there is an analogy between various physical quantities in rotational and translational motion.
 - ii) In the rotational motion about a fixed axis, the moment of inertia (I) is the analogy the mass is the linear motion is the measure of inertia of the body.
 - iii) Therefore moment of inertia is also regarded as the translatory motion.
 - iv) In case of translatory motion the inertia of the body depends totally upon its mass but in case of rotational motion the moment of inertia of the body not only depends on the mass of the body but also on distribution of mass about the given axis of rotation.
 - v) The analogy between various physical quantities in two types of motion.



Sr. No.	Translational Motion.	Rotational Motion.
1.	Mass = m	Moment of inertia = I
2.	Displacement = s at x	Angular displacement = θ
3.	Velocity = v	Angular velocity
4.	Acceleration = a	Angular Acceleration = α
5.	Force = $F = ma$	Torque = $T = I\alpha$
6.	Linear momentum	Angular momentum $L = I\omega$
7.	Kinetic energy $\frac{1}{2} mv^2$	Rotational K.E = $\frac{1}{2} I\omega^2$
8.	Work done = $F \cdot s$	Work done = $T \cdot \theta$

2) Derive an expression for M.I of a solid cylinder its axis of symmetry.

1) Let M be the mass, R be the radius and h be length of solid cylinder.





- 2) The mass per unit length of the solid cylinder M/L . Let yy' be the axis passing through O and perpendicular to its own axis xx' as

Fig.

- 3) To find its moment of inertia imagine that the cylinder is made up of large number thin discs i.e. thickness of the disc is dx . Let us consider one of such disc at a distance x from O . The thickness of the disc is dx and its mass is $(M/L) dx$.

- 4) Therefore, moment of inertia of the disc about its diameter mass dm is $= \frac{(\text{radius})^2}{4}$

$$= \frac{m}{L} dx \cdot \frac{R^2}{4}$$

- 5) According to the principle of parallel axis a moment of inertia of the disc about the axis $yy' = \frac{m}{L} dx \cdot \frac{R^2}{4} + \frac{m}{L} dx \cdot x^2$

- 6) The moment of inertia of the whole cylinder about the axis yy' can be obtained by integrating the above equation between the limits $x=0$ to $x=L/2$



and multiplying the result by 2.
 7) Therefore the moment of inertia of the cylinder about the axis yy' is.

$$I = \int dI = 2 \int_0^{L/2} \left(\frac{m}{L} \frac{R^4}{4} dx + \frac{m}{L} x^2 dx \right)$$

$$= \frac{2m}{L} \left[\frac{R^4}{4} x + \frac{x^3}{3} \right]_0^{L/2}$$

$$= \frac{2m}{L} \left[\frac{R^4}{4} \cdot \frac{L}{2} + \frac{L^3}{8 \times 3} \right]$$

$$= \frac{2m}{L} \left[\frac{R^4 L}{8} + \frac{L^3}{24} \right]$$

$$I = m \left[\frac{R^4}{4} + \frac{L^2}{12} \right]$$

This is the required expression.



Name - Yagita Yasant Sutar

Class - BSC III

Sem - I

Roll NO - 8206

13

INTERNAL EXAM

VIVEKANAND COLLEGE, KOLAPUR

Answer the following question.

1) Explain analogy of rotational motion with translational motion.

→ i) It is noticed that there is an analogy between various physical quantities in rotational and translational motion.

ii) In rotation motion about a fixed axis the moment of inertia (I) is an analogy to the mass (m) in linear motion. But mass is the inertia of the body.

iii) therefore moment of inertia is also regarded as the rotational inertia.

iv) The moment of inertia of the body in case of rotational motion plays the same role as the mass of the body in translatory motion.

v) In case of translatory motion the inertia of the body depend totally upon its mass but in case of rotational motion the moment of inertia depends on the mass of the



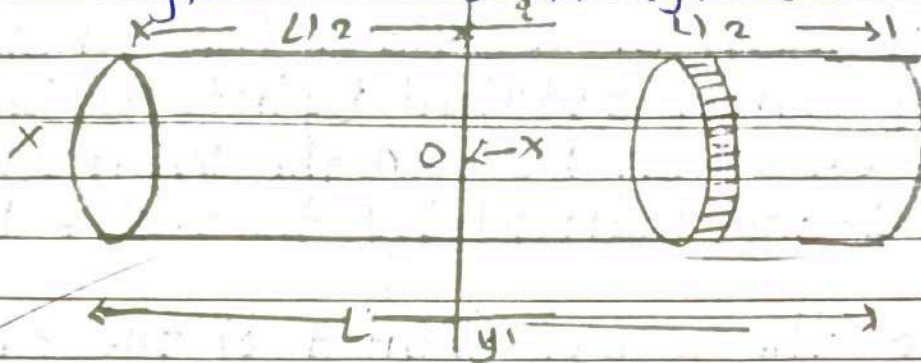
mass of the body but also an distribution of mass about the given axis of rotation
 vii) The analogy between various physical quantities in two type of motion.

Sr. No	Translational Motion	Rotational Motion
1.	Mass = m	moment of inertia = I
2.	displacement = s or x	Angular displacement
3.	velocity = v	Angular Velocity = ω
4.	Acceleration = a	Angular Accertion = α
5.	Force = $F = ma$	Torque = $\tau = I\alpha$
6.	linear momentum = $p = mv$	Angular momentum = $L = I\omega$
7.	Kinetic energy $\frac{1}{2} mv^2$	Rotation K.E = $\frac{1}{2} I\omega^2$
8.	Work done = $F \cdot s$	Work done = $\tau \cdot \theta$



2) Derive an expression for M.I. of a solid cylinder about its axis symmetry.

→ i) let M be the mass R be the radius and L be the length of the solid cylinder.



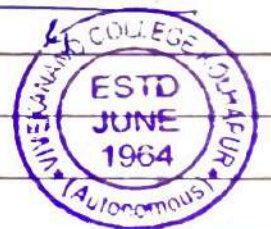
2) the mass of per unit length of the solid cylinder is m/L let Y be the axis passing through its centre and perpendicular to its own axis $x-x'$ as shown in

Fig

3) To find its moment of inertia imagine that the cylinder to be made up of large number of thin discs let us consider one of such disc at distance x from a the thickness of the disc is dx its mass is $(M/L) dx$.

4) therefore moment of inertia of the disc about its diameter = mass of disc \times $(\text{radius})^2$

$$= \frac{m}{L} dx \times \frac{R^2}{2}$$



5) According to the principle of parallel axis a.

moment of inertia of the disc about the axis yy' = $\frac{m}{l} dz$.

$$\frac{R^2}{4} + \frac{m}{l} dz \cdot z^2$$

6) The moment of inertia of the whole cylinder the yy' can be obtained by integration the and above equation between the limits $x=0$ to $l/2$ and multiplying the result by 2.

7) therefore the moment of inertia of the cylinder about the axis yy' is

$$I = \int dI = 2 \int_0^{l/2} (m/l \cdot R^2/4 dz + m/l \cdot z^2 dz)$$

$$= 2 \frac{m}{l} \left[\frac{R^2}{4} z + \frac{z^3}{3} \right]_{0}^{l/2}$$

$$= 2 \frac{m}{l} \left[\frac{R^2}{4} \cdot \frac{l}{2} + \frac{l^3}{3 \times 3} \right]$$

$$= 2 \frac{m}{l} \left[\frac{R^2 \cdot l}{8} + \frac{l^3}{24} \right]$$

$$I = m \left[\frac{R^2}{4} + \frac{l^2}{12} \right]$$

This is the required expression



Name - Madhavi Dhondiram Shingare. Department - Physics.
Class - Bsc. III Div - B. Roll NO - 7859 8203
Sem - V INTERNAL EXAM

VIVEKANAND COLLEGE, KOLAPUR

Q. 1. Write the correct alternatives.

- 1) Mass is the measure of inertia in linear motion.
- 2) Acceleration of a body rolling down in an inclined plane is independent of mass of the body.
- 3) Force in rotational motion is analogous to torque in translational motion.
- 4) Moment of inertia of a spherical shell about its diameter is $\frac{2}{3} mR^2$.

Q. 2. Answer the following questions.

1) Explain analogy of rotational motion with translational motion.

-
- i) It is noticed that there is an analogy between various physical quantities in rotational and translational motion.
 - ii) In rotational motion about a fixed axis, the moment of inertia (I) is analogous to the mass (m) in linear motion. But mass in linear motion is the measure of inertia of the body.
 - iii) Therefore moment of inertia is also regarded as the rotational inertia.
 - iv) The moment of inertia of the body in case of rotational motion plays the same role as the mass of the body in translatory motion.
 - v) In case of translatory motion, the inertia of the body depends totally upon its mass, but in case of rotational motion the moment of inertia of the body not only depends on the mass of the body but also on its distribution of mass about the given axis.



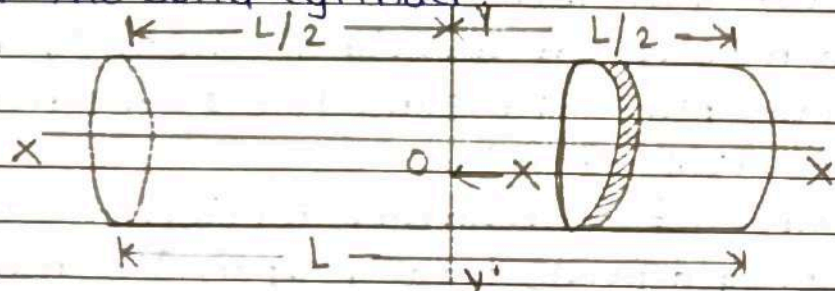
of rotation.

vi) The analogy between various physical quantities in two types of motion

Sl. No.	Translational Motion	Rotational Motion
1.	Mass = m	Moment of Inertia = I
2.	Displacement = s or x	Angular displacement = θ
3.	Velocity = v	Angular velocity = ω
4.	Acceleration = a	Angular Acceleration = α
5.	Force = $F = ma$	Torque = $\tau = I\alpha$
6.	Linear Momentum = $P = mv$	Angular momentum = $L = I\omega$
7.	Kinetic energy = $\frac{1}{2}mv^2$	Rotational K.E = $\frac{1}{2}I\omega^2$
8.	Work done = $F \cdot s$	Work done = $\tau \cdot \theta$

2) Derive an expression for M.I of a solid cylinder about its axis of symmetry.

1) Let M be the mass, R be the radius and L be the length of the solid cylinder.



2) The mass per unit length of the solid cylinder is M/L . Let YY' be the axis passing through its centre O and perpendicular to its own axis XX' , as shown in

fig.

3) To find its moment of inertia, imagine that the cylinder to be made up of large number of thin discs. Let us consider one of such disc at a distance x from O . The thickness of the disc is dx , obviously, mass of the disc is $(M/L) dx$.

4) Therefore, moment of inertia of the disc about its diameter = mass of disc \times $\frac{(\text{radius})^2}{4}$

$$= \frac{m}{L} dx \cdot \frac{R^2}{4}$$

4) According to the principle of parallel axis, a moment of inertia of the disc about the axis $YY' = \frac{m}{L} dx \cdot \frac{R^2}{4} + \frac{m}{L} dx \cdot x^2$

5) The moment of inertia of the whole cylinder about the axis yy' can be obtained by integrating the above equation between the limits $x = 0$ to $L/2$ and multiplying the result by 2.

6) Therefore, the moment of inertia of the cylinder about the axis yy' is,

$$I = \int dI = 2 \int_0^{L/2} \left(\frac{m}{L} \frac{R^4}{4} dx + \frac{m}{L} x^2 dx \right)$$

$$= 2 \frac{m}{L} \left[\frac{R^4}{4} x + \frac{x^3}{3} \right]_0^{L/2}$$

$$= 2 \frac{m}{L} \left[\frac{R^4}{4} \cdot \frac{L}{2} + \frac{L^3}{8 \times 3} \right]$$

$$= 2 \frac{m}{L} \left[\frac{R^4 \cdot L}{8} + \frac{L^3}{24} \right]$$

$$I = m \left[\frac{R^4}{4} + \frac{L^2}{12} \right]$$

This is the required expression.



Shri Swami Vivekanand Shikshan Sanstha's
Vivekanand College, Kolhapur (Autonomous)

Internal Examination 2022-23

PHYSICS-DSC -1001E

B.Sc. – III, Sem – V Quantum Mechanics

Time: 30 Minutes

Marks: 20

Q. 1. LONG Answer Questions (Any one)

(10)

- i) Obtain Schrodinger , s time independent equation and time dependent equation
 - ii) Using cartesian components of operators L_x , L_y & L_z prove that $[L_x , L_y] = i \hbar L_z$ &
- $[L^2 , L_x] = 0$

Q. 2. SHORT Answer Questions (any two)

(10)

- i) Show that $[x , P_x] = i \hbar$ give its physical significance
- ii) Give physical significance of wave function
- iii) Obtain Schrodingers equation in spherical polar coordinate system for hydrogen atom



॥ ज्ञान, विज्ञान आणि सुसंस्कार यांसाठी शिक्षण प्रसार ॥

- शिक्षणमहर्षी डॉ. बापूजी साळुंखे

34605

Shri Swami Vivekanand Shikshan Sanstha Kolhapur's

VIVEKANAND COLLEGE, KOLHAPUR (AUTONOMOUS)

SUPPLIMENT

16
20

Suppliment No. :

Roll No. : 8214

Class : B.Sc III, Sem-V

Signature
of
Supervisor

Subject : Quantum Mechanics

Test / Tutorial No. :

Div. :

Q.2)

$$1) [x, p_x] \psi = [x, -i\hbar \frac{\partial}{\partial x}] \psi$$

$$= -i\hbar [x, \frac{\partial}{\partial x}] \psi$$

$$= -i\hbar \left(x \cdot \frac{\partial}{\partial x} - \frac{\partial}{\partial x} \cdot x \right) \psi$$

$$= -i\hbar \left\{ x \cdot \frac{\partial \psi}{\partial x} - \frac{\partial}{\partial x} (x\psi) \right\}$$

$$= -i\hbar \left\{ x \cdot \frac{\partial \psi}{\partial x} - x \frac{\partial \psi}{\partial x} - \psi \right\}$$

$$= -i\hbar (-\psi)$$

$$[x, p_x] \psi = i\hbar \psi$$

$$[x, p_x] = i\hbar$$



2) Physical Significance of wave function -

- i) Ψ must be finite for all values of x, y, z .
- ii) Ψ must be single valued i.e. for each set of values of x, y, z , Ψ must have only one value.
- iii) Ψ must be continuous in all regions, except in those regions where the potential energy $V(x, y, z) = \infty$
- iv) Ψ must vanish at infinity i.e. $\Psi = 0$ as $x \rightarrow \pm\infty$, $y \rightarrow \pm\infty$ or $z \rightarrow \pm\infty$.

v) The potential derivatives of Ψ i.e. $\frac{\partial \Psi}{\partial x}$, $\frac{\partial \Psi}{\partial y}$, $\frac{\partial \Psi}{\partial z}$ must also be finite, single-valued and continuous at all points except at points where the potential $V(x)$ is infinite.



Q. 17 2)

(i) The orbital angular momentum is defined as,

$$L = r \times p$$

where, p represents momentum of body and r is the position vector of a body from the axis of rotation.

The components of momentum x, y, z in cartesian co-ordinates will be

$$\left. \begin{aligned} L_x &= -i\hbar \left(y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right) \\ L_y &= -i\hbar \left(z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right) \\ L_z &= -i\hbar \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right) \end{aligned} \right\} \text{--- (1)}$$

Let us write the commutation relation $[L_x, L_y]$ as,

$$[L_x, L_y] = L_x L_y - L_y L_x$$

Substituting operator values we get,

$$L_x L_y = -i\hbar \left(y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right) (-i\hbar) \left(z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right)$$

$$= -\hbar^2 \left(y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right) \left(z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right)$$

$$= -\hbar^2 \left[y \frac{\partial}{\partial z} \left(z \frac{\partial}{\partial x} \right) - y \frac{\partial}{\partial z} \left(x \frac{\partial}{\partial z} \right) - z \frac{\partial}{\partial y} \left(z \frac{\partial}{\partial x} \right) + \right.$$

$$\left. z \frac{\partial}{\partial y} \left(x \frac{\partial}{\partial z} \right) \right]$$

$$= -\hbar^2 \left[yz \frac{\partial^2}{\partial z \partial x} + 0 + y \frac{\partial}{\partial z} - yx \frac{\partial^2}{\partial z^2} - 0 - z^2 \frac{\partial^2}{\partial y \partial x} \right.$$

$$\left. - 0 + \frac{zx \partial^2}{\partial y \partial z} + 0 \right]$$

$$= -\hbar^2 \left\{ \left(z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right) \left(y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right) \right\}$$



$$= -\hbar^2 \left\{ z \frac{\partial}{\partial x} \left(y \frac{\partial}{\partial z} \right) - z \frac{\partial}{\partial z} \left(z \frac{\partial}{\partial y} \right) - x \frac{\partial}{\partial z} \left(y \frac{\partial}{\partial z} \right) + x \frac{\partial}{\partial z} \left(z \frac{\partial}{\partial y} \right) \right\}$$

$$= -\hbar^2 \left\{ zy \frac{\partial^2}{\partial x \partial z} + 0 - \frac{z^2 \partial^2}{\partial x \partial y} - 0 - xy \frac{\partial^2}{\partial z^2} - 0 + xz \frac{\partial^2}{\partial z \partial y} + x \frac{\partial}{\partial y} \right\}$$

$$= -\hbar^2 \left\{ zy \frac{\partial^2}{\partial x \partial z} - z^2 \frac{\partial^2}{\partial x \partial y} - xy \frac{\partial^2}{\partial z^2} + xz \frac{\partial^2}{\partial z \partial y} + x \frac{\partial}{\partial y} \right\}$$

On subtracting above two relations we get,

$$[L_x L_y] = \hbar^2 \left\{ \left(yz \frac{\partial^2}{\partial z \partial x} + y \frac{\partial}{\partial x} - yx \frac{\partial^2}{\partial z^2} - z^2 \frac{\partial^2}{\partial y \partial x} + zx \frac{\partial^2}{\partial y \partial z} \right) - \left(zy \frac{\partial^2}{\partial x \partial z} - z^2 \frac{\partial^2}{\partial x \partial y} + xy \frac{\partial^2}{\partial z^2} + xz \frac{\partial^2}{\partial z \partial y} + x \frac{\partial}{\partial y} \right) \right\}$$

$$= -\hbar^2 \left(y \frac{\partial}{\partial x} - x \frac{\partial}{\partial y} \right)$$

$$= -\hbar^2 \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right)$$

$$= i\hbar(-i\hbar) \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right)$$

$$[L_x L_y] = i\hbar L_z$$



Shri Swami Vivekanand Shikshan Sanstha Kolhapur's

VIVEKANAND COLLEGE, KOLHAPUR (AUTONOMOUS)**SUPLIMENT**Signature
of
Supervisor

Subject : Quantum Mechanics

Test / Tutorial No. :

Div. :

Suppliment No. :

Roll No. : 8201

Class : BSc-III, Sem-V

18
20

Q.1}

2) ① The orbital angular momentum is defined as,

$$L = r \times p$$

where, p represents momentum of body and r is the position vector of a body from the axis of rotation.

The components of momentum x, y, z in cartesian co-ordinates will be.

$$L_x = -i\hbar \left(y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right) \quad \text{--- ①}$$

$$L_y = -i\hbar \left(z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right)$$

$$L_z = -i\hbar \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right)$$

Let us write the commutation relation $[L_x, L_y]$ as

$$[L_x, L_y] = L_x L_y - L_y L_x$$

Substituting operator values we get

$$\begin{aligned} L_x L_y &= -i\hbar \left(y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right) (-i\hbar) \left(z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right) \\ &= -\hbar^2 \left(y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right) \left(z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right) \end{aligned}$$



$$= -\hbar^2 \left[y \frac{\partial}{\partial z} \left(z \frac{\partial}{\partial x} \right) - y \frac{\partial}{\partial z} \left(x \frac{\partial}{\partial z} \right) - z \frac{\partial}{\partial y} \left(z \frac{\partial}{\partial x} \right) + z \frac{\partial}{\partial y} \left(x \frac{\partial}{\partial z} \right) \right]$$

$$= -\hbar^2 \left[yz \frac{\partial^2}{\partial z \partial x} + 0 + y \frac{\partial}{\partial x} - yx \frac{\partial^2}{\partial z^2} - 0 - \frac{z^2 \partial^2}{\partial y \partial x} - 0 + zx \frac{\partial^2}{\partial y \partial z} + 0 \right]$$

$$= -\hbar^2 \left\{ \left(z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right) \left(y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right) \right\}$$

$$= -\hbar^2 \left\{ z \frac{\partial}{\partial x} \left(y \frac{\partial}{\partial z} \right) - z \frac{\partial}{\partial x} \left(z \frac{\partial}{\partial y} \right) - x \frac{\partial}{\partial z} \left(y \frac{\partial}{\partial z} \right) + x \frac{\partial}{\partial z} \left(z \frac{\partial}{\partial y} \right) \right\}$$

$$= -\hbar^2 \left\{ zy \frac{\partial^2}{\partial x \partial z} + 0 - \frac{z^2 \partial^2}{\partial x \partial y} - 0 - xy \frac{\partial^2}{\partial z^2} - 0 + xz \frac{\partial^2}{\partial z \partial y} + x \frac{\partial}{\partial y} \right\}$$

$$= \hbar^2 \left\{ zy \frac{\partial^2}{\partial x \partial z} - \frac{z^2 \partial^2}{\partial x \partial y} - xy \frac{\partial^2}{\partial z^2} + xz \frac{\partial^2}{\partial z \partial y} + x \frac{\partial}{\partial y} \right\}$$

On Subtracting above two relations we get,

$$[L_x, L_y] = \hbar^2 \left\{ \left(yz \frac{\partial^2}{\partial z \partial x} + y \frac{\partial}{\partial x} + yx \frac{\partial^2}{\partial z^2} - \frac{z^2 \partial^2}{\partial y \partial x} + zx \frac{\partial^2}{\partial y \partial z} \right) - \left(zy \frac{\partial^2}{\partial x \partial z} - \frac{z^2 \partial^2}{\partial x \partial y} - xy \frac{\partial^2}{\partial z^2} + xz \frac{\partial^2}{\partial z \partial y} + x \frac{\partial}{\partial y} \right) \right\}$$



$$= -\hbar^2 \left(y \frac{\partial}{\partial x} - x \frac{\partial}{\partial y} \right)$$

$$= \hbar^2 \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right)$$

$$= i\hbar(-i\hbar) \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right)$$

$$[L_x, L_y] = i\hbar L_z$$

(ii) The total angular momentum is defined by the relation
$$L^2 = L_x^2 + L_y^2 + L_z^2$$

Let us take,

$$[L^2, L_x] = [L_x^2 + L_y^2 + L_z^2, L_x]$$

$$= [L_x^2 + L_y^2 + L_z^2] L_x - L_x [L_x^2 + L_y^2 + L_z^2]$$

$$= L_x^2 L_x + L_y^2 L_x + L_z^2 L_x - (L_x L_x^2 + L_x L_y^2 + L_x L_z^2)$$

$$= (L_y^2 L_x - L_x L_y^2) + (L_z^2 L_x - L_x L_z^2)$$

$$= [L_y^2, L_x] + [L_z^2, L_x]$$

We know that,

$$[ab, c] = a[b, c] + [a, c]b$$

$$[L^2, L_x] = [L_y L_y, L_x] + [L_z L_z, L_x]$$

$$= L_y [L_y, L_x] + [L_y, L_x] L_y + L_z [L_z, L_x] + [L_z, L_x] L_z$$

$$= (-i\hbar L_z) L_y - (i\hbar L_z) L_y + i\hbar L_y L_z + i\hbar L_y L_z$$

$$[L^2, L_x] = 0$$



$$\begin{aligned}
 Q.27) [x, p_x] \psi &= [x, -i\hbar \frac{\partial}{\partial x}] \psi \\
 &= -i\hbar [x, \frac{\partial}{\partial x}] \psi \\
 &= -i\hbar \left(x \frac{\partial}{\partial x} - \frac{\partial}{\partial x} \cdot x \right) \psi \\
 &= -i\hbar \left\{ x \frac{\partial \psi}{\partial x} - \frac{\partial}{\partial x} (x\psi) \right\} \\
 &= -i\hbar \left\{ x \frac{\partial \psi}{\partial x} - x \frac{\partial \psi}{\partial x} - \psi \right\} \\
 &= -i\hbar (-\psi)
 \end{aligned}$$

$$[x, p_z] \psi = i\hbar \psi$$

$$[x, p_z] = i\hbar$$

Q.27) Physical significance of wave function -

- 1) ψ must be finite for all values of x, y, z
- 2) ψ must be single valued i.e. for each set of values of x, y, z , ψ must have only one value.
- 3) ψ must be continuous in all regions, except in those regions where the potential energy $V(x, y, z) = \infty$
- 4) ψ must vanish at infinity i.e. $\psi = 0$ as $x \rightarrow \pm\infty$, $y \rightarrow \pm\infty$ or $z \rightarrow \pm\infty$.

5) The partial derivatives of ψ , i.e. $\frac{\partial \psi}{\partial x}$, $\frac{\partial \psi}{\partial y}$, $\frac{\partial \psi}{\partial z}$ must also be finite, single-valued and continuous at all points except at points where the potential $V(x)$ is infinite.



Shri Swami Vivekanand Shikshan Sanstha's
Vivekanand College, Kollhapur (Autonomous)

Internal Examination 2022-23

PHYSICS-DSC -1001F

B.Sc. – III, Sem – V Nuclear and particle Physics

Marks: 20

Time: 30 Minutes

Q. 1. LONG Answer Questions (Any one)

(10)

1. Define specific heat at constant volume and specific heat at constant pressure. Obtain an expression for $C_P - C_V$. Apply that for perfect gas and for van-derwaal's gas.
2. Derive Planck's law of radiation in terms of frequency and wavelength.

Q. 2. SHORT Answer Questions (any two)

(10)

1. Write a note on perfect black body?
2. What is electron gas? Obtain expression for fermi energy of electrons.
3. Explain concept of energy density.



Shri Swami Vivekanand Shikshan Sanstha Kolhapur's

VIVEKANAND COLLEGE, KOLHAPUR (AUTONOMOUS)**SUPPLIMENT**Signature
of
SupervisorSubject: Nucleos and particle
physics

Test / Tutorial No.:

Div.:

Suppliment No. :

20
20

Roll No. : 8210

Class : B.Sc III, Sem V

Q2.

1 Perfect black body

→ A very good experimental approximation of black body is provided by cavity the interior walls of which are maintained at uniform temperature and which communicates with outside through a hole having very small diameter in comparison with dimensions of the cavity. Any radiation entering the hole is partly absorbed and partly diffused reflected with only negligible fraction eventually finding its way out of the hole. This is true regardless of material of which interior walls are composed. Such black body is called Ferry's black body.

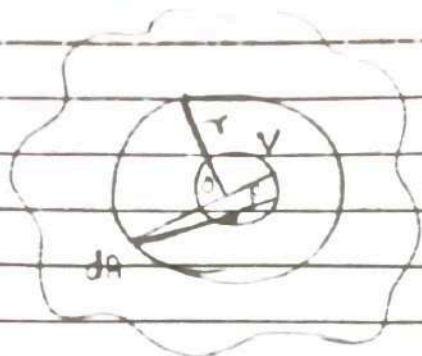
The radiant energy emitted within the cavity whose walls are at temperature θ is equal to radiant emittance of black body at the same temperature.

For this reason, the radiation within the cavity is called black body radiation.

black body radiation is homogeneous and isotropic. The physical laws for perfect gas can, as they be applied to the black body radiation.



3. Concept of Energy Density



Energy density of diffuse radiation.

The energy contained in unit volume of radiation is called energy density of radiation. Let us find the energy density of diffused radiation inside uniformly heated enclosure density inside enclosure at large distance from its walls. All the radiation contained in V may be assumed to come from sphere is very large as compared to the dimensions of volume V .

The amount of radiation energy contained in V due to this elementary cone is

$$k \frac{dA \cdot f \cdot t}{r^2} = k \cdot \frac{dA}{r^2} \frac{f l}{c}$$

The total amount of radiation energy coming from dA and contained in V is found by summing over all

$$\int k \frac{dA \cdot f l}{r^2} = k \frac{dA}{c} \int \frac{f l}{r^2} = k \frac{dA}{c} \int \frac{V}{r^2}$$

$$E f l = \text{Volume element } V$$

Q1.

1. Planck applied quantum theory to the problem of black body radiation. The radiation from black body is supposed to consist of photons of energy, ranging from zero to practically infinite. They are supposed to be moving in all possible directions with speed of light c . The wall of enclosure can be replaced by emission of several photons of frequencies $\nu_1, \nu_2, \nu_3, \dots$ so that total energy of the system is constant.

$$h\nu = h\nu_1 + h\nu_2 + h\nu_3 + \dots$$

The total number of eigen states between the momentum range P and $P + dP$ is

$$n \cdot g(P) dP = \frac{4\pi P^2 dP}{h^3/v} = \frac{4\pi v P^3}{h^3} dP$$

$$\text{photon, } P = \frac{h\nu}{c} \text{ and } dP = \frac{h d\nu}{c}$$

put eqⁿ we have the number of given states between the frequency range ν to $\nu + d\nu$ is,

$$g(\nu) d\nu = \frac{4\pi v \nu^2}{c^3} d\nu$$

As there are two modes of propagation for each photon, the total no. of eigen states available for the photons in the frequency range ν and $\nu + d\nu$ is,

$$g(\nu) d\nu = \frac{8\pi v \nu^2}{c^3} d\nu$$

Consider n_i Bose-Einstein distribution law

$$n_i = \frac{g_i}{e^{h\nu_i/kT} - 1}$$



put $\alpha = 0$ and $h\nu = h\nu$ we have number of photons with frequency between ν and $\nu + d\nu$ in volume V of radiation

$$N(\nu)d\nu = \frac{8\pi V \nu^2}{c^3(e^{h\nu/KT} - 1)} d\nu$$

$$\therefore n(\nu)d\nu = \frac{N(\nu)d\nu}{V} = \frac{8\pi d\nu}{c^3(e^{h\nu/KT} - 1)}$$

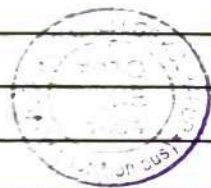
Hence the energy density $u(\nu)d\nu$ which is the energy per unit volume in frequency range $\nu + \nu + d\nu$ is

$$u(\nu)d\nu = h\nu \times n(\nu)d\nu$$

$$\therefore u(\nu)d\nu = \frac{8\pi h \nu^3}{c^3} \frac{d\nu}{(e^{h\nu/KT} - 1)}$$

Eqⁿ represent plank - radiation formula in terms of the frequency of black - body radⁿ

10



Shri Swami Vivekanand Shikshan Sanstha Kolhapur's

VIVEKANAND COLLEGE, KOLHAPUR (AUTONOMOUS)**SUPPLIMENT**17
20

Suppliment No. :

Roll No. : 8217

Class : B.Sc.-III, Sem.-IV

Signature
of
Supervisor

Subject : Nuclear and Particle Physics.

Test / Tutorial No. :

Div. :

Q. 12

2) Plank applied quantum theory to the problem of black body radiation. The radiation from a black body is supposed to consist of photons of energy, ranging from zero to practically infinite. They are supposed to be moving in all possible directions with the speed of light c . The wall of the enclosure, it can be replaced by the emission of several photons of frequencies $\nu_1, \nu_2, \nu_3, \dots$ so that the total energy of the system is constant.

$$h\nu = h\nu_1 + h\nu_2 + h\nu_3 + \dots$$

Thus, in this case the principle of conservation of number of particles is not valid i.e. $\sum \delta n_i \neq 0$. It is equivalent to say that the undetermined constant α is equal to zero. Moreover, the photons are Bose particles with spin 1 having two modes of propagation. At any instant all photons having their momenta between $P + dP$ will lie within a spherical shell with radius P and thickness dP . The volume of this shell is $4\pi P^2 dP$.

Therefore, the total number of eigen states between the momentum range, P and $P + dP$ is

$$g(p) dp = \frac{4\pi P^2 dP}{h^3/V} = \frac{4\pi V P^3}{h^3} dp$$

for a photon, $p = \frac{h\nu}{c}$ and $dp = \frac{h d\nu}{c}$



Substituting eqⁿ we have the number of given states between the frequency range ν and $\nu + d\nu$ is

$$g(\nu) d\nu = \frac{4\pi V \nu^2}{c^3} d\nu$$

As there are two modes of propagation for each photon the total number of eigen states available for the photons in the frequency range ν and $\nu + d\nu$ is

$$g(\nu) d\nu = \frac{8\pi V \nu^2}{c^3} d\nu$$

Now, consider the Bose-Einstein distribution law.

$$n_i = \frac{g_i}{e^{\alpha} e^{u_i/kT} - 1}$$

Putting $\alpha = 0$ and $u_i = h\nu$, we have the number of photons with frequency between ν and $\nu + d\nu$ in the volume V of the radiation.

$$N(\nu) d\nu = \frac{8\pi V \nu^2}{e^3 (e^{h\nu/kT} - 1)} d\nu$$

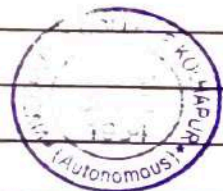
Therefore, the number of photons with frequency in the range ν and $\nu + d\nu$ per unit volume of the radiation is,

$$n(\nu) d\nu = \frac{N(\nu) d\nu}{V} = \frac{8\pi d\nu}{c^3 (e^{h\nu/kT} - 1)}$$

Hence, the energy density $u(\nu) d\nu$, which is the energy per unit volume in the frequency range ν and $\nu + d\nu$ is,

$$u(\nu) d\nu = h\nu \times n(\nu) d\nu$$

$$\therefore u(\nu) d\nu = \frac{8\pi h \nu^3}{c^3} \frac{d\nu}{(e^{h\nu/kT} - 1)}$$



Eqⁿ represents Planck-radiation formula in terms of the frequency of black-body radiation.

Q22

3) Energy Density of radiation \rightarrow

The energy contained in unit volume of radiation is called the energy density of radiation. Let us find the energy density of diffused radiation inside a uniformly heated enclosure of any shape. Consider a very small element of volume inside the enclosure at large distance from its walls. All the radiation contained in V may be assumed to come from a sphere described about any point O inside V . The radius r of the sphere is very large as compared to the dimensions of volume V .

The amount of radiation energy contained in V due to this elementary cone is,

$$k \frac{dA \cdot \frac{V}{r^2}}{r^2} = k_1 \frac{dA}{r^2} \frac{V}{r}$$

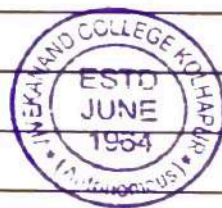
The total amount of radiation energy coming from dA and contained in V is found by summing over all elementary cones, as,

$$\sum k \frac{dA}{r^2} \cdot \frac{V}{r} = \frac{k}{c} \frac{dA}{r^2} \epsilon f l = \frac{k}{c} \frac{dA}{r^2} V$$

where, $\epsilon f l =$ Volume element V .

1) Black body \rightarrow

A very good experimental approximation of a black body is provided by a cavity the interior walls of which are maintained at a uniform temperature and which with the outside through a hole having very small diameter in comparison with the dimensions of the cavity. Any radiation entering the hole is partly absorbed and partly different reflected with only a negligible fraction eventually finding its way out of the hole.



Shri Swami Vivekanand Shikshan Sanstha's
Vivekanand College, Kolhapur (Autonomous)

Internal Examination 2022-23

PHYSICS-DSC -1001F

B.Sc. – III, Sem – V Mathematical Physics

Time: 30 Minutes

Marks: 20

Q. 1. LONG Answer Questions (Any one)

(10)

1. Obtain an expression for curl of vector field in orthogonal curvilinear co-ordinators.
2. What is mean by an ensemble. Discuss microcanonical and canonical ensemble.

Q. 2. SHORT Answer Questions (any two)

(10)

1. Write a note on assessable microstate.
2. What is electron gas? Obtain expression for fermi energy of electrons.
3. Discuss cylindrical co-ordinate system.



॥ ज्ञान, विज्ञान आणि सुसंस्कार यासाठी शिक्षण प्रसार ॥

- शिक्षणमहर्षी डॉ. बापूजी साळुंखे

34604

Shri Swami Vivekanand Shikshan Sanstha Kolhapur's

VIVEKANAND COLLEGE, KOLHAPUR (AUTONOMOUS)

SUPPLIMENT

Signature
of
Supervisor

Suppliment No. :

Roll No. : 8216

Class : B.Sc. III, Sem V

Subject : Mathematical physics

Test / Tutorial No. : Internal exam

Div. :

15
20

Q1

2. The collection of large number of assemblies is called as an ensemble. All the member of an ensemble are referred as elements. These elements differ in their microscopic states. Thus an ensemble is defined as collection of large number of assemblies which are essentially independent of one another but which have been made macroscopically as identical as possible.

a) Microcanonical Ensemble -

The microcanonical ensemble is collection of essentially independent assemblies having same Energy E , volume V , number N of system all the system are of same type. The individual assemblies are separated by rigid impermeable and well insulated wall such that the values of E , V and N , are not affected by presence of other systems, we can't actually specify the macroscopic energy of an assembly exactly.



E, V, N	E, V, N	E, V, N
E, V, N	E, V, N	E, V, N
E, V, N	E, V, N	E, V, N

Due to complete isolation of system, thermodynamics can't be applied to this ensemble. Microcanonical ensemble can be made practical use if we can apply thermodynamics to it. This is possible only average energy \bar{E} of the system is specified

b. Canonical ensemble -

The Canonical ensemble is collection of essentially independent assemblies having same temperature T , volume V and the no. of identical particles N .

To assure that all the assemblies in thermal contact with each other. The individual assemblies are separated by rigid, impermeable but diathermic walls, since energy can be exchanged between the assemblies and reach common temperature, thus in Canonical ensemble, system can change but not particles.

T, V, N	T, V, N	T, V, N
T, V, N	T, V, N	T, V, N
T, V, N	T, V, N	T, V, N



Thermodynamics can be applied to such ensemble. In thermodynamics we do not know the exact value of energy as we usually deal with systems kept in thermal contact.

Q2.

1. Note On Microstates.

Consider an assembly consisting of large number of independent system / number of molecules in phase space. The state of individual / molecules may be separated in phase space by point known as phase point / representative point. The phase space may be divided into cells $1, 2, 3 \dots i$. A phase point for any of the system / molecules may reside in one of these cells.

Microstate is an arrangement of specified system / molecules with their representative point in particular cell. In other words, microstate of assembly may be defined by the specifications of individual position of phase point for each system or molecule of assemblies.

A phase point gives the state of motion of the molecule at that point, so to define the microstate of the assembly we should know the state of each and every molecule of assembly a given instant of time.



॥ ज्ञान, विज्ञान आणि सुसंस्कार यांसाठी शिक्षण प्रसार ॥

- शिक्षणमहर्षी डॉ. बापूजी साळुंखे

34603

Shri Swami Vivekanand Shikshan Sanstha Kolhapur's

VIVEKANAND COLLEGE, KOLHAPUR (AUTONOMOUS)

SUPPLIMENT

Signature
of
Supervisor

Suppliment No. :

Roll No. : 8212

Class : B.Sc III

Subject : Mathematical physics

Test / Tutorial No. : Internal exam

Div. :

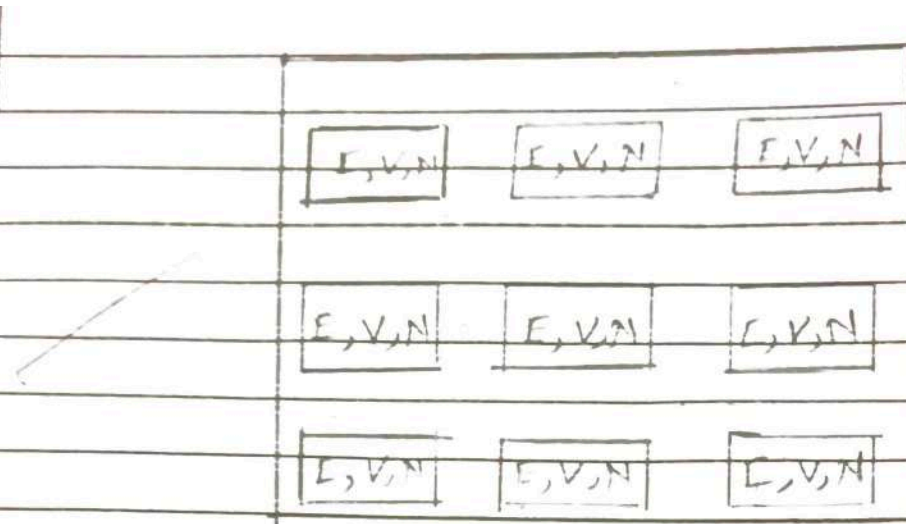
Q.1

2. The collection of large number of assemblies is called as an ensemble. All the member of an ensemble are referred as elements. These elements differ in their microscopic states. Thus an ensemble is defined as a collection of a large number of assemblies which are essentially independent of one another but which have been made macroscopically as identical as possible.

a) Microcanonical ensemble' -

The microcanonical ensemble is a collection of essentially independent assemblies having the same energy E , volume V and number N of system, all the systems are of the same type. The individual assemblies are separated by rigid, impermeable, and well insulated wall such that the values of E , V and N are not affected by the presence of other systems. We cannot actually specify the macroscopic energy of an assembly exactly.





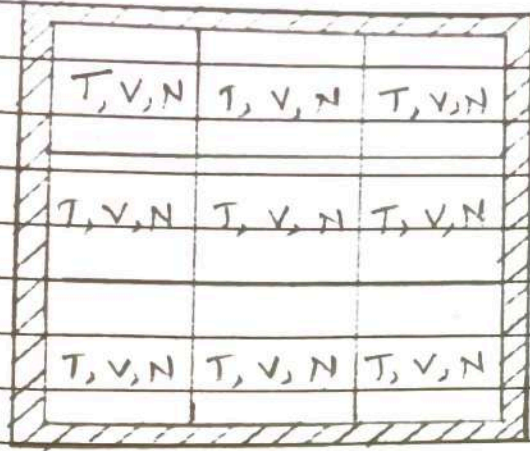
Due to the complete isolation of the system, thermodynamics cannot be applied to this ensemble. Microcanonical ensemble can be made of practical use if we can apply thermodynamics to it. This is possible only the average energy E of the system is specified.

b) Canonical ensemble :-

The canonical ensemble is collection of essentially independent assemblies having the same temperature T , volume V and the number of identical particles N .

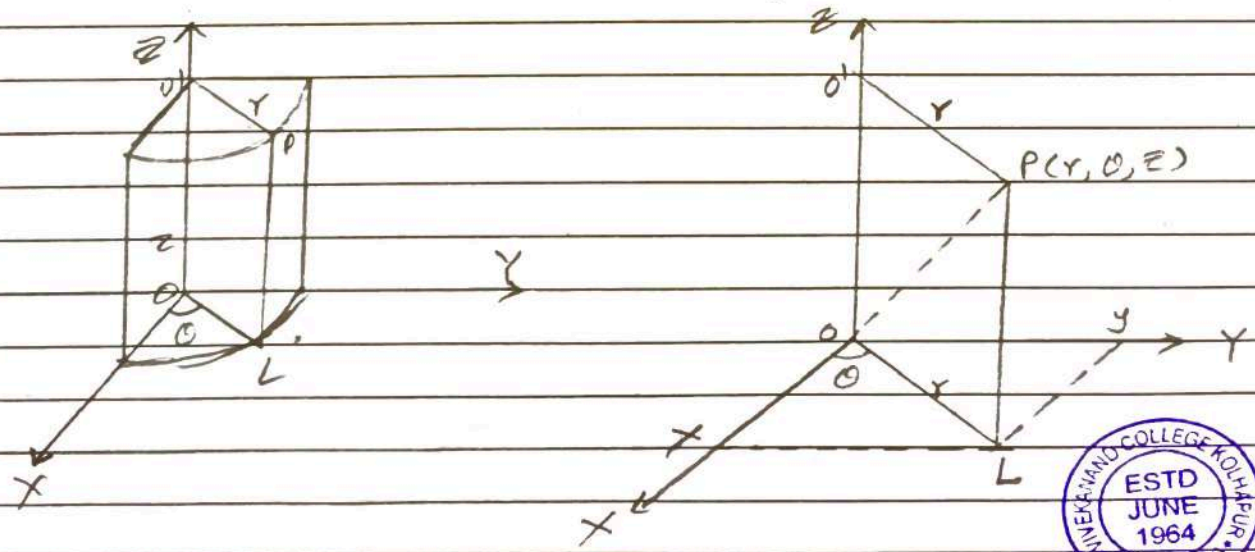
To assure that all the assemblies have the same temperature we could simply bring all the assemblies in thermal contact with each other. Fig. represents symbolically a canonical ensemble. The individual assemblies are separated by rigid, impermeable but diathermic walls, since energy can be exchanged between the assemblies and will reach a common temperature. Thus in canonical ensemble, system can exchange energy but not particles.





Q.2

3. In cylindrical co-ordinate system, there is symmetry about an axis, such as OZ as shown in fig. cartesian co-ordinates of point P in space are (x, y, z) . In order to find cylindrical co-ordinates of point P , we draw a radius vector r from O' as shown in fig., then we draw Pl parallel to OZ and then $O'l$ parallel to $O'P$. Then angle lOx thus formed is θ . Then position of point in cylindrical co-ordinate is represented by (r, θ, z)



So from fig. we have

$$\left. \begin{aligned} x &= r \cos \theta \\ y &= r \sin \theta \\ z &= z \end{aligned} \right\}$$

Differentiating we get

$$dx = dr \cos \theta - r \sin \theta d\theta$$

$$dy = dr \sin \theta + r \cos \theta d\theta$$

$$dz = dz$$

Then expression for line segment in cartesian co-ordinate is

$$ds^2 = dx^2 + dy^2 + dz^2$$

substituting values of dx , dy and dz , we get

$$ds^2 = dr^2 + r^2 d\theta^2 + dz^2$$

