

Vivekanand College, Kolhapur. (Autonomous)
Department of Physics
Internal Examination Notice
2018-19

Date:30/09/2018

All students of class B.Sc. I, B.Sc. II and B.Sc. III are hereby noticed that the first term internal evaluation examination is scheduled as per following time table.

Nature of question paper:

For B.Sc. I : Long answer question (Any one from given two questions) for 10 marks

Short answer question (Any two from given three questions) for 10 marks

For B.Sc. II : Long answer question (Any one from given two questions) for 10 marks

Short answer question (Any two from given three questions) for 10 marks

For B.Sc. II (Astro) : Long answer question (Any one from given two questions) for 10 marks

Short answer question (Any two from given three questions) for 10 marks

For B.Sc. III : Long answer question (Any one from given two questions) for 10 marks

Short answer question (Any two from given three questions) for 10 marks


Internal Evaluation Examination 2018-19.

SEM I, SEM III and SEM V

Time Table

Sr. No.	Class	Paper	Date	Time
1.	B.Sc. I	Paper I	11/10/2018	11:00 am to 12:00 pm
2.	B.Sc. II	Paper III	11/10/2018	11:00 am to 12:00 pm
3.	B.Sc. II (Astrophysics)	Paper I	12/10/2018	11:00 am to 12:00 pm
4.	B.Sc. III	Paper V (section I)	15/10/2018	11:00 am to 12:00 pm
		Paper V (section II)		01:00 am to 2:00 pm
		Paper VI (section I)	16/10/2018	11:00 am to 12:00 pm
		Paper VI (section II)		01:00 am to 2:00 pm




 HOD
 Head of the
 Department of Physics
 Vivekanand College, Kolhapur

Vivekananda College Kolhapur (Autonomous).
Department of Physics: Internal examination 2018-19
B.Sc. III Semester V
Subject: Atomic and Molecular Spectra, Astronomy and
Astrophysics

Marks: 20 (Each question carry one mark)

Time : 20 min

Q.1 Attempt any ONE

(10)

1. Discuss the principle of proton-synchrotron with a special reference to two step acceleration.
2. Explain the principle of electron-synchrotron with special reference to two-step acceleration.

Q.2 Attempt any TWO

(10)

1. Discuss different methods used to measure nuclear radius.
2. What are nucleons? Explain their intrinsic properties.
3. What is the shape and size of nucleus?



Shri Swami Vivekanand Shikshan Sanstha's

Vivekanand College, Kolhapur

(Autonomous)

Department of Physics

Internal exam

B.Sc. III Sem V

Attendance Sheet

Roll No.	Name Of The Student	Signature			
		15-10-2018	15-10-2018	16-10-2018	16-10-2018
8501	Aniket Nandkumar Chile	Achile	Achile	Achile	Achile
8502	Shubham Nandkumar Chodankar	Chodankar	Chodankar	Chodankar	Chodankar
8503	Ankita Jayawant Chougule	A	A	A	A
8504	Patil Pramod Dashrath	Ppatil	Ppatil	Ppatil	Ppatil
8505	Ankita Ravindra Digraje	D	D	D	D
8506	Pooja Lagamana Ghulanawar	Gpooja	Gpooja	Gpooja	Gpooja
8507	Prasad Rajaram Gulavani	Gulavani	Gulavani	Gulavani	Gulavani
8508	Vinayak Baburao Kesarkar	VC	VC	VC	VC
8509	Aishwarya Sanjay Kumbhar	Sanjay	Sanjay	Sanjay	Sanjay
8510	Karale Prajakta Mansing	Karale	Karale	Karale	Karale
8511	Shamal Vijay Mohite	SM	SM	SM	SM
8512	Tejaswini Tanaji Musale	Musale	Musale	Musale	Musale
8513	Anisa Ajj Nadaf	Anisa	Anisa	Anisa	Anisa
8514	Somesh Krishnat Nerlekar	Nerlekar	Nerlekar	Nerlekar	Nerlekar
8515	Sourabh Sanjay Patil	Patil	Patil	Patil	Patil
8516	Anuja Uday Patil	Apatil	Apatil	Apatil	Apatil
8517	Pranil Yuvraj Patil	Patil	Patil	Patil	Patil
8518	Pratiraj Sampat Patil	Patil	Patil	Patil	Patil
8519	Satish Shivaji Patil	Patil	Patil	Patil	Patil
8520	Sheral Shivaji Patil	Patil	Patil	Patil	Patil
8521	Shrinath Dhondiram Shinde	Shinde	Shinde	Shinde	Shinde
8522	Kumbhar Swaroop Sunil	K	K	K	K
8523	Ajit Sadashiv Thorat	Thorat	Thorat	Thorat	Thorat
8524	Ruhan Eliyas Ustad	Ustad	Ustad	Ustad	Ustad
8525	Vaibhav Vasant Yadav	Vyadav	Vyadav	Vyadav	Vyadav

Internal Examiner.....

[Handwritten Signature]



VIVEKANAND COLLEGE, KOLHAPUR (AUTONOMOUS)

SUPPLIMENT

Signature of Supervisor

Suppliment No. :

Roll No. : 8516

Class : B.Sc III, Sem-

Subject : Atomic & molecular Physics

Test / Tutorial No. : Internal Exam

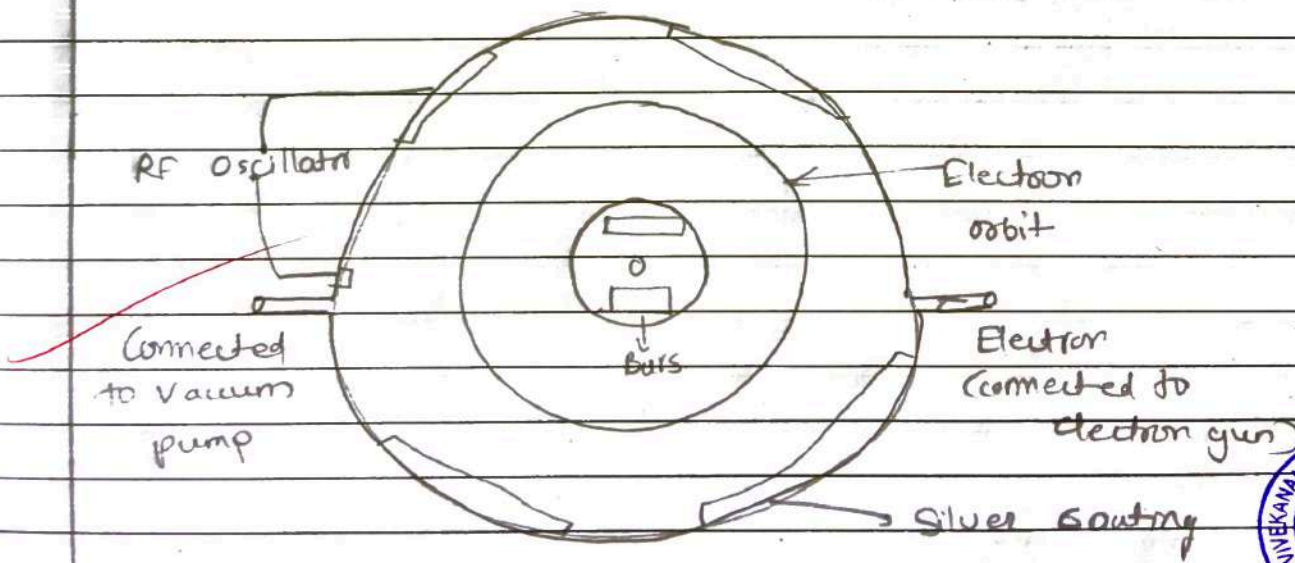
Div. :

10
20

Q.15

2> Electron - Synchrotron

Construction- Synchrotron also uses doughnut-shaped vacuum chamber in AC magnetic field. The weight of the magnet is reduced in synchrotron. The vacuum chamber inside c-shaped magnet. The magnetic focusing is required, the pole faces are constructed to provide maximum field at centre.



Working -

With the help of electron gun, electrons are injected into the vacuum chamber with energy range of up to 100 keV. After the electrons are accelerated to high energy, then the electrons may attain velocity comparable to that of light. The bars play a very important role. Once the steel bars get saturated, they no longer obey Faraday's law of electromagnetic induction. The electrons gain energy after every revolution. The magnetic field and orbit radius decide the energy of electrons. After the electrons gain maximum energy, radio-frequency oscillator is turned off and larger current is sent through auxiliary coils so that the electrons change their orbit radius due to unstable magnetic field. The target decreases highly energetic X-rays when electrons strike the target. The electrons can attain energy up to 330 MeV by synchrotron action and 7.5 MeV by betatron action.



Shri Swami Vivekanand Shikshan Sanstha Kolhapur's

VIVEKANAND COLLEGE, KOLHAPUR (AUTONOMOUS)

SUPPLIMENT

Signature
of
Supervisor

Subject : (Atomic and molecular physics)

Test / Tutorial No. : Internal

Div. :

Suppliment No. :

Roll No. : 8522

Class : B.Sc III

Q.1

1. Proton synchrotron can be used to accelerate neutrons, alpha particles in addition to protons.

construction:-

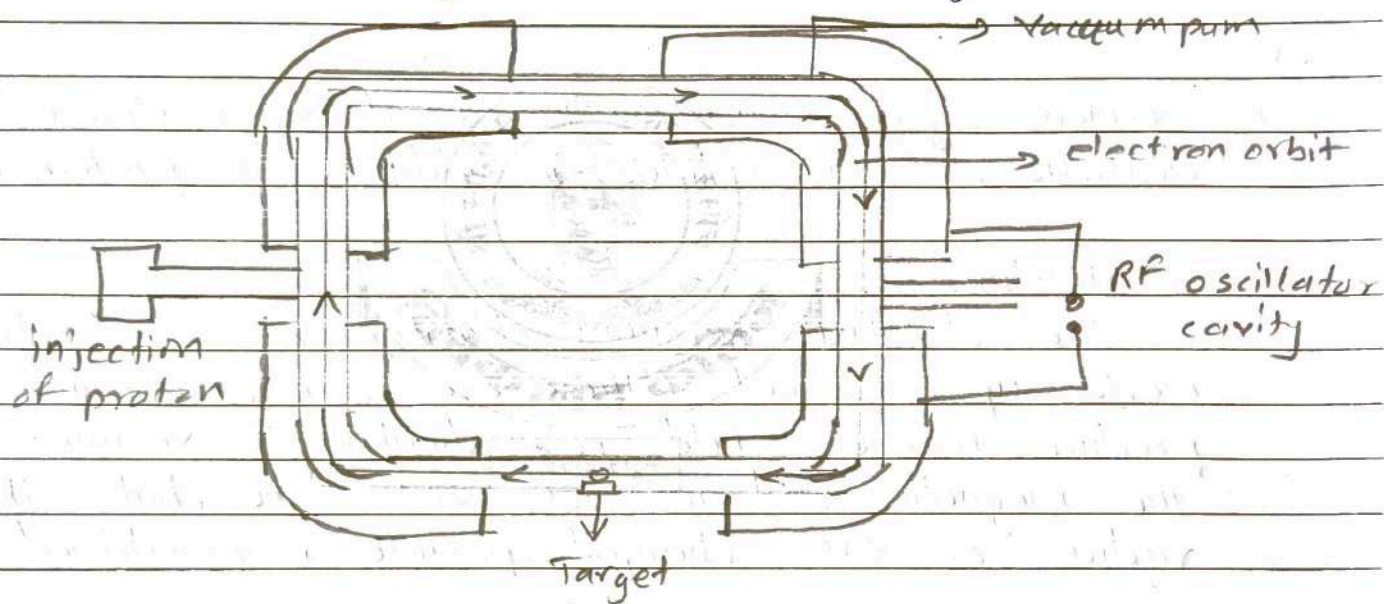
The doughnut vacuum chamber of synchrotron is made up of steel. There are four quadrants that produce magnetic field perpendicular to vacuum chamber. The magnetic field increases with time, but the radius of "+ve" charged particle is maintained constant.

working:-

Using a linear accelerator such as Van de Graff generator, the protons are accelerated towards the doughnut chamber hence, initially the "+ve" particle such as proton can be accelerated up to 10 MeV. These particles are injected when the magnetic field is small. These electrons then come under the influence of radio frequency oscillator. Moreover, the magnetic field



These electrons then also increased to keep the electrons in circular orbit frequency constant at constant radius. As proton completes its revolution, it almost gains an impulse of 1 kV/turn which increases its energy as well as frequency. Then to maintain the phase stability the freq. of radiofreq. oscillator is also increased in order to synchronize it with the freq. of proton. The range up to which protons can be accelerated is higher than the range of electrons.



When the protons are accelerated to maximum energy levels, then the radio-freq. is distorted, so that the radius of orbit changes. After the proton gets out of its track, it will strike the target.

frequency of revolution of positive particles to be

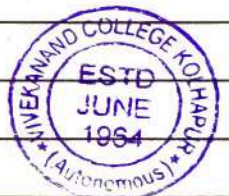
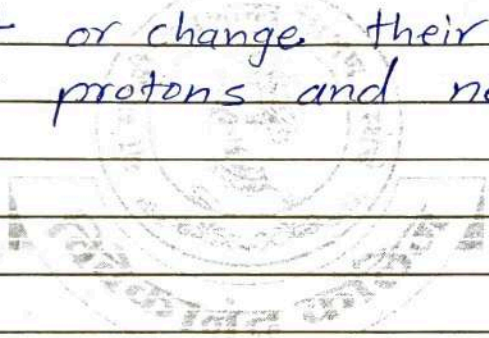
$$f = \frac{(q B c^2) (2\pi r)}{2\pi (k + m_0 c^2) (2\pi r_0 + 4d)}$$



Q.2

3.

The nucleus makes up much less than 0.01% of the volume of the atom, but typically contains more than 99.9% of the mass of the atom. The chemical properties of substance are determined by the negatively charged electrons enshrouding the nucleus. Most nuclei are spherical or ellipsoidal, though some exotic shapes exist. Nuclei can vibrate and rotate when struck by other particles. Some are unstable and will break apart or change their relative number of protons and neutrons.



VIVEKANAND COLLEGE, KOLHAPUR (AUTONOMOUS)

SUPPLIMENT

Suppliment No. :

Roll No. : 8525

Class : B.Sc-III

Signature
of
Supervisor

Subject : Atomic and molecules

Test / Tutorial No. : Internal

Div. :

- Q.1. proton Synchrotron Can be used to accelerate
2. deuterons, alpha particles, in addition to protons

Construction -

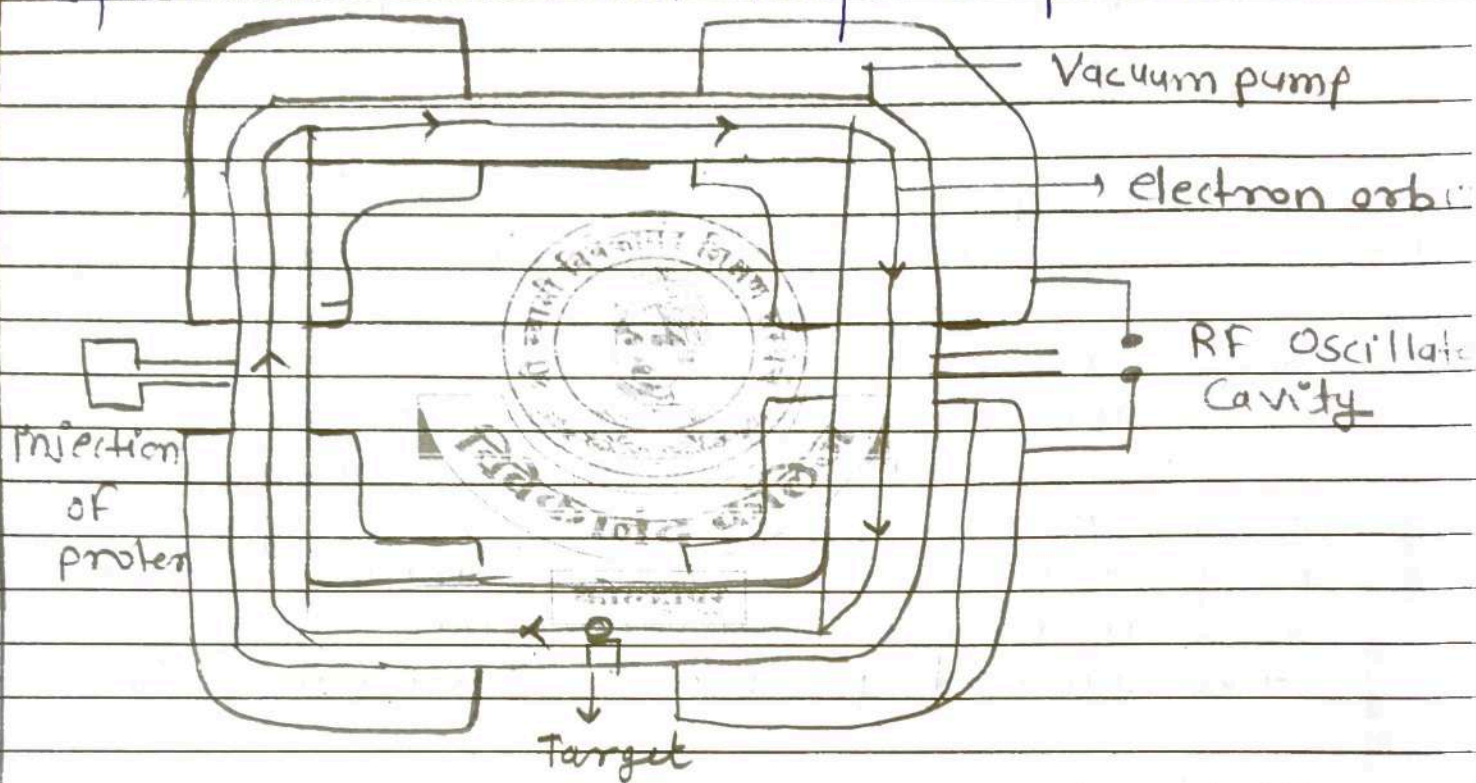
The doughnut vacuum chamber of Synchrotron is made up of steel. There are four quadrants that produce magnetic field perpendicular to vacuum chamber. The magnetic field increases with time, but the radius of +ve charged particle is maintained constant

Working -

Using linear accelerator such as Van de graff generator, the protons are accelerated towards the doughnut chamber, hence initially the +ve particle such as proton can be accelerated up to 10 MeV. These particles are injected when the magnetic field is small. These electrons then come under the influence of Oscillator.



Also increased to keep the electrons in circular orbit constant radius. As proton completes its revolution, it almost gains an impulse of $\pm kv/\text{turn}$ which increases its energy as well as frequency. Hence to maintain the phase stability the frequency of radiofrequency oscillator is also increased in order to synchronise it with freq of proton.



when the protons are accelerated to maximum energy levels then radio-frequency is distorted, so that the radius of orbit changes. After the proton gets out of its track, it will strike.

frequency of revolution of positive particle to be

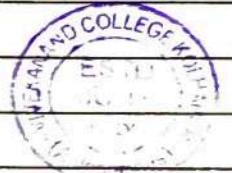


$$F = \frac{(9Be^2)(2\pi E)}{2\pi(K + m_0c^2)(2\pi r_0 + 4e)}$$

Q2.

3.

The nucleus makes up much less than 0.01 of the volume of atom, but typically contains more than 99.9% of the mass of the atom. The chemical properties of substance are determined by the negativity charged e^-



Vivekananda College Kolhapur (Autonomous).
Department of Physics: Internal examination 2018-19

B.Sc. III Semester V
Subject: Mathematical Physics

Marks: 20 (Each question carry one mark)

Time : 20 min

Q.1 Attempt any ONE (10)

1. Discuss Hamilton variational principle.
2. Derive Hamilton's canonical equation of motion from variational principle.

Q.2 Attempt any TWO (10)

1. State equivalence of Lagrange's and Newton's equations.
2. Write a note on degree of freedom and constraints.
3. What is relation between H and L?



॥ ज्ञान, विज्ञान आणि सुसंस्कार यांसाठी शिक्षण प्रसार ॥

- शिक्षणमहर्षी डॉ. बापूजी साळुंके

34025

Shri Swami Vivekanand Shikshan Sanstha Kolhapur's

VIVEKANAND COLLEGE, KOLHAPUR (AUTONOMOUS)

SUPPLIMENT

Signature
of
Supervisor

Suppliment No. :

Roll No. : 8517

Class : BSc-III, Sem-V

Subject : Mathematical & Statistical
Physics

Test / Tutorial No. :

Div. :

20
20

Q.17

1) The curl of a vector in a orthogonal curvilinear Co-ordinate system

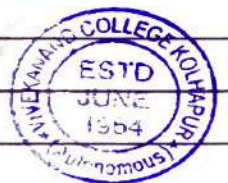
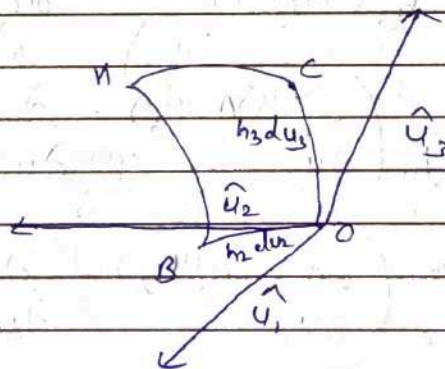
$$\oint_S (\nabla \times \vec{A}) \cdot d\vec{s} = \oint \vec{A} \cdot d\vec{l}$$

Using mean value theorem for integral

$$\text{curl } \vec{A} \cdot \int_S d\vec{s} = \oint \vec{A} \cdot d\vec{l}$$

$$\oint \vec{A} \cdot d\vec{l} = \oint \vec{A} \cdot d\vec{l} + \oint \vec{A} \cdot d\vec{l} + \oint \vec{A} \cdot d\vec{l} + \oint \vec{A} \cdot d\vec{l}$$

let $\vec{A} = A_1 \vec{u}_1 + A_2 \vec{u}_2 + A_3 \vec{u}_3$



$$\oint \vec{A} \cdot d\vec{l} = A_2 h_2 du_2 + (A_3 h_3 du_3 + \frac{\partial A_3 h_3}{\partial u_2} du_2 du_3) + (-A_2 h_2 du_2 - \frac{\partial A_2 h_2}{\partial u_3} du_2 du_3) + (-A_3 h_3 du_3)$$

$$\oint \vec{A} \cdot d\vec{l} = \left[\frac{\partial A_3 h_3}{\partial u_2} - \frac{\partial A_2 h_2}{\partial u_3} \right] du_2 du_3$$

$$(\nabla \times \vec{A})_{u_1} = \frac{\left[\frac{\partial A_3}{\partial u_2} - \frac{\partial A_2}{\partial u_3} \right] h_2 h_3 du_2 du_3}{h_2 h_3 du_2 du_3}$$

$$= \frac{1}{h_1 h_3} \left(\frac{\partial A_3 h_3}{\partial u_2} - \frac{\partial A_2 h_2}{\partial u_3} \right)$$

similarly by choosing areas S_2 & S_3

$$(\nabla \times \vec{A})_{u_2} = \frac{1}{h_3 h_1} \left(\frac{\partial (A_1 h_1)}{\partial u_3} - \frac{\partial (A_3 h_3)}{\partial u_1} \right)$$

$$(\nabla \times \vec{A})_{u_3} = \frac{1}{h_1 h_2} \left(\frac{\partial (A_2 h_2)}{\partial u_1} - \frac{\partial (A_1 h_1)}{\partial u_2} \right)$$

$$\text{curl } \vec{A} = \nabla \times \vec{A} = \hat{u}_1 \left(\frac{\partial (A_3 h_3)}{\partial u_2} - \frac{\partial (A_2 h_2)}{\partial u_3} \right)$$

$$+ \frac{\hat{u}_2}{h_3 h_1} \left(\frac{\partial (A_3 h_3)}{\partial u_3} - \frac{\partial (A_2 h_2)}{\partial u_1} \right)$$

$$+ \frac{u_3}{h_1 h_2} \left(\frac{\partial (A_2 h_2)}{\partial u_1} - \frac{\partial (A_1 h_1)}{\partial u_2} \right)$$



$$\text{curl } \vec{A} = \vec{\nabla} \times \vec{A} = \frac{1}{h_1 h_2 h_3} \begin{vmatrix} u_1 h_1 & u_2 h_2 & u_3 h_3 \\ \frac{\partial}{\partial u_1} & \frac{\partial}{\partial u_2} & \frac{\partial}{\partial u_3} \\ A_1 h_1 & A_2 h_2 & A_3 h_3 \end{vmatrix}$$

Q2) Gradient

let us consider, $\phi(u_1, u_2, u_3)$ a scalar function.

$$(\text{grad } \phi)_{u_1} = (\nabla \phi)_{u_1} = \lim_{\delta u_1 \rightarrow 0} \frac{\phi(B) - \phi(A)}{AB}$$

where $\phi(B)$ and $\phi(A)$ are the values of scalar function ϕ at B and A separated by distance AB δu_1 .

$$AB = h_1 \delta u_1 \quad (\text{grad } \phi)_{u_1} = (\nabla \phi)_{u_1} = \lim_{\delta u_1 \rightarrow 0} \left(\frac{\delta \phi}{h_1 \delta u_1} \right)$$

$$(\nabla \phi)_{u_1} = \frac{1}{h_1} \frac{\partial \phi}{\partial u_1}$$

similarly, the component of gradient of ϕ in the direction u_2 & u_3

$$(\text{grad } \phi)_{u_2} = (\nabla \phi)_{u_2} = \frac{1}{h_2} \frac{\partial \phi}{\partial u_2}$$

$$(\text{grad } \phi)_{u_3} = (\nabla \phi)_{u_3} = \frac{1}{h_3} \frac{\partial \phi}{\partial u_3}$$

If $\hat{u}_1, \hat{u}_2, \hat{u}_3$ are unit

$$\text{grad } \phi = \nabla \phi = \frac{u_1}{h_1} \frac{\partial \phi}{\partial u_1} + \frac{u_2}{h_2} \frac{\partial \phi}{\partial u_2} + \frac{u_3}{h_3} \frac{\partial \phi}{\partial u_3}$$

$$\nabla \phi = \left(\frac{\hat{u}_1}{h_1} \frac{\partial}{\partial u_1} + \frac{\hat{u}_2}{h_2} \frac{\partial}{\partial u_2} + \frac{\hat{u}_3}{h_3} \frac{\partial}{\partial u_3} \right) \phi$$

It gives the operator $(\vec{\nabla})$ in orthogonal curvilinear

$$\vec{\nabla} = \frac{\hat{u}_1}{h_1} \frac{\partial}{\partial u_1} + \frac{\hat{u}_2}{h_2} \frac{\partial}{\partial u_2} + \frac{\hat{u}_3}{h_3} \frac{\partial}{\partial u_3}$$



32 Laplacian operator ∇^2 in orthogonal curvilinear

we know that,

$$\operatorname{div} \vec{A} = \vec{\nabla} \cdot \vec{A} = \frac{1}{h_1 h_2 h_3} \left(\frac{\partial}{\partial u_1} (\vec{A}_1 h_2 h_3) + \frac{\partial}{\partial u_2} (\vec{A}_2 h_1 h_3) + \frac{\partial}{\partial u_3} (\vec{A}_3 h_1 h_2) \right)$$

let $\vec{A} = \vec{\nabla} \cdot f$, where f is scalar field substituting we

$$\operatorname{div} \vec{A} = \vec{\nabla} \cdot \vec{\nabla} f = \nabla^2 f$$

$$= \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u_1} (h_2 h_3 \vec{\nabla} f_{u_1}) + \frac{\partial}{\partial u_2} (h_3 h_1 \nabla f_{u_2}) + \frac{\partial}{\partial u_3} (h_1 h_2 \nabla f_{u_3}) \right]$$

$$\nabla^2 f = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u_1} \left(\frac{h_2 h_3}{h_1} \frac{\partial f}{\partial u_1} \right) + \frac{\partial}{\partial u_2} \left(\frac{h_1 h_3}{h_2} \frac{\partial f}{\partial u_2} \right) + \frac{\partial}{\partial u_3} \left(\frac{h_1 h_2}{h_3} \frac{\partial f}{\partial u_3} \right) \right]$$

The Laplacian operator ∇^2 in orthogonal curvilinear coordinate is

$$\nabla^2 = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u_1} \left(\frac{h_2 h_3}{h_1} \frac{\partial}{\partial u_1} \right) + \frac{\partial}{\partial u_2} \left(\frac{h_1 h_3}{h_2} \frac{\partial}{\partial u_2} \right) + \frac{\partial}{\partial u_3} \left(\frac{h_1 h_2}{h_3} \frac{\partial}{\partial u_3} \right) \right]$$



VIVEKANAND COLLEGE, KOLHAPUR (AUTONOMOUS)

SUPPLIMENT

15
20

Suppliment No. :

Roll No. : 8519

Class : BSc. III Sem - V

Signature
of
Supervisor

Subject : Mathematical & statistical
Physics

Test / Tutorial No. :

Div. :

Laplacean operator -

We know that

$$\text{div. } \vec{A} = \vec{\nabla} \cdot \vec{A} = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u_1} (A_1 h_2 h_3) + \frac{\partial}{\partial u_2} (A_2 h_3 h_1) + \frac{\partial}{\partial u_3} (A_3 h_1 h_2) \right]$$

Let $\vec{A} = \vec{\nabla} \cdot f$, where f is scalar field

substitute we get,

$$\text{div. } \vec{A} = \vec{\nabla} \cdot \vec{\nabla} f = \nabla^2 f$$

$$= \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u_1} (h_2 h_3 (\vec{\nabla} f)_{u_1}) + \frac{\partial}{\partial u_2} (h_3 h_1 (\vec{\nabla} f)_{u_2}) + \frac{\partial}{\partial u_3} (h_1 h_2 (\vec{\nabla} f)_{u_3}) \right]$$

$$\nabla^2 f = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u_1} \left(\frac{h_2 h_3}{h_1} \frac{\partial f}{\partial u_1} \right) + \frac{\partial}{\partial u_2} \left(\frac{h_1 h_3}{h_2} \frac{\partial f}{\partial u_2} \right) + \frac{\partial}{\partial u_3} \left(\frac{h_1 h_2}{h_3} \frac{\partial f}{\partial u_3} \right) \right]$$

The Laplacian operator ∇^2 in orthogonal curvilinear co-ordinate is given by,

$$\nabla^2 = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u_1} \left(\frac{h_2 h_3}{h_1} \frac{\partial}{\partial u_1} \right) + \frac{\partial}{\partial u_2} \left(\frac{h_1 h_3}{h_2} \frac{\partial}{\partial u_2} \right) + \frac{\partial}{\partial u_3} \left(\frac{h_1 h_2}{h_3} \frac{\partial}{\partial u_3} \right) \right]$$



Curl of Vector in OCCS -

The curl of a vector may be obtained by using Stoke's theorem which is given by,

$$\iint_S (\nabla \times \vec{A}) ds = \oint \vec{A} dl$$

Using mean value theorem for integral

$$\text{Curl } \vec{A} \iint_S ds = \oint \vec{A} dl \quad \text{or} \quad \text{Curl } \vec{A} = \lim_{ds \rightarrow 0} \frac{\oint \vec{A} dl}{\iint_S ds}$$

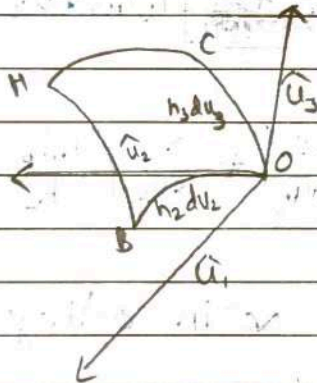
To evaluate this consider surface (S) OBHC of a parallelepiped normal to \hat{u}_1 at O. Refer Fig. 1.6. Boundary of surface denotes closed line.

Then closed line integral can be written as

$$\oint_{OBHC} \vec{A} dl = \oint_{OB} \vec{A} dl + \oint_{BH} \vec{A} dl + \oint_{HC} \vec{A} dl + \oint_{CO} \vec{A} dl$$

Let,

$$\vec{A} = A_1 \hat{u}_1 + A_2 \hat{u}_2 + A_3 \hat{u}_3$$



Evaluating integral, we get

$$\begin{aligned} \oint_{OBHC} \vec{A} dl &= A_2 h_2 du_2 + \left(A_3 h_3 du_3 + \frac{\partial A_3 h_3}{\partial u_2} du_2 du_3 \right) \\ &+ \left(-A_2 h_2 du_2 - \frac{\partial A_2 h_2}{\partial u_3} du_2 du_3 \right) + (-A_3 h_3 du_3) \end{aligned}$$

$$\therefore \oint_{OBHC} \vec{A} dl = \left[\frac{\partial A_3 h_3}{\partial u_2} - \frac{\partial A_2 h_2}{\partial u_3} \right] du_2 du_3$$



Dividing by area of S_1 (OBHC) equal to $h_2 h_3 du_2 du_3$ and taking limit du_2 & du_3 approaches to zero.

$$(\nabla \times \vec{A})_{u_1} = \frac{\left[\frac{\partial A_3 h_3}{\partial u_2} - \frac{\partial A_2 h_2}{\partial u_3} \right] du_2 du_3}{h_2 h_3 du_2 du_3}$$

$$= \frac{1}{h_1 h_3} \left(\frac{\partial A_3 h_3}{\partial u_2} - \frac{\partial A_2 h_2}{\partial u_3} \right)$$

Similarly by choosing areas S_2 & S_3 perpendicular to \hat{u}_2 and \hat{u}_3 respectively.

$$(\nabla \times \vec{A})_{u_2} = \frac{1}{h_3 h_1} \left(\frac{\partial (A_1 h_1)}{\partial u_3} - \frac{\partial (A_3 h_3)}{\partial u_1} \right)$$

$$(\nabla \times \vec{A})_{u_3} = \frac{1}{h_1 h_2} \left(\frac{\partial (A_2 h_2)}{\partial u_1} - \frac{\partial (A_1 h_1)}{\partial u_2} \right)$$

From these three components resultant eqⁿ is,

$$\text{curl } \vec{A} = \nabla \times \vec{A} = \frac{\hat{u}_1}{h_2 h_3} \left(\frac{\partial (A_3 h_3)}{\partial u_2} - \frac{\partial (A_2 h_2)}{\partial u_3} \right)$$

$$+ \frac{\hat{u}_2}{h_3 h_1} \left(\frac{\partial (A_1 h_1)}{\partial u_3} - \frac{\partial (A_3 h_3)}{\partial u_1} \right)$$

$$+ \frac{\hat{u}_3}{h_1 h_2} \left(\frac{\partial (A_2 h_2)}{\partial u_1} - \frac{\partial (A_1 h_1)}{\partial u_2} \right)$$

$$\text{curl } \vec{A} = \nabla \times \vec{A} = \frac{1}{h_1 h_2 h_3} \begin{vmatrix} \hat{u}_1 h_1 & \hat{u}_2 h_2 & \hat{u}_3 h_3 \\ \frac{\partial}{\partial u_1} & \frac{\partial}{\partial u_2} & \frac{\partial}{\partial u_3} \\ A_1 h_1 & A_2 h_2 & A_3 h_3 \end{vmatrix}$$

VIVEKANAND COLLEGE, KOLHAPUR (AUTONOMOUS)

SUPLIMENT

17
20

Suppliment No. :

Roll No. : 8516

Class : BSc. III Sem V

Signature
of
Supervisor

Subject : Mathematical & statistical Physics

Test / Tutorial No. :

Div. :

Gradient of scalar field in OCCS -

Let us consider $\phi(u_1, u_2, u_3)$ a scalar function. The gradient of the scalar field ϕ in direction of u_1 -axis can be written as.

$$(\text{grad } \phi)_{u_1} = (\nabla \phi)_{u_1} = \lim_{\delta u_1 \rightarrow 0} \frac{\phi(B) - \phi(A)}{AB}$$

where $\phi(B)$ and $\phi(A)$ are values of scalar function ϕ at B and A separated by distance $AB = h_1 \delta u_1$. The quantity $[\phi(B) - \phi(A)]$ may be taken as increase in ϕ on travelling distance $AB = h_1 \delta u_1$.

This may be written as $\delta \phi$ for the limiting case where $\delta u_1 \rightarrow 0$

$$(\text{grad } \phi)_{u_1} = (\nabla \phi)_{u_1} = \lim_{\delta u_1 \rightarrow 0} \left[\frac{\delta \phi}{h_1 \delta u_1} \right]$$

$$(\nabla \phi)_{u_1} = \frac{1}{h_1} \frac{\partial \phi}{\partial u_1}$$

Similarly, the component of gradient of ϕ in the direction u_2 and u_3 axes are.

$$(\text{grad } \phi)_{u_2} = (\nabla \phi)_{u_2} = \frac{1}{h_2} \frac{\partial \phi}{\partial u_2}$$

$$(\text{grad } \phi)_{u_3} = (\nabla \phi)_{u_3} = \frac{1}{h_3} \frac{\partial \phi}{\partial u_3}$$

If $\hat{u}_1, \hat{u}_2, \hat{u}_3$ are unit vectors along the u_1, u_2, u_3 directions respectively then we can write,



$$\text{grad } \phi = \nabla \phi = \frac{\hat{u}_1}{h_1} \frac{\partial \phi}{\partial u_1} + \frac{\hat{u}_2}{h_2} \frac{\partial \phi}{\partial u_2} + \frac{\hat{u}_3}{h_3} \frac{\partial \phi}{\partial u_3}$$

or

$$\nabla \phi = \left[\frac{\hat{u}_1}{h_1} \frac{\partial}{\partial u_1} + \frac{\hat{u}_2}{h_2} \frac{\partial}{\partial u_2} + \frac{\hat{u}_3}{h_3} \frac{\partial}{\partial u_3} \right] \phi$$

It gives the del operator (∇) in orthogonal curvilinear co-ordinates.

$$\nabla = \frac{\hat{u}_1}{h_1} \frac{\partial}{\partial u_1} + \frac{\hat{u}_2}{h_2} \frac{\partial}{\partial u_2} + \frac{\hat{u}_3}{h_3} \frac{\partial}{\partial u_3}$$

Laplacean - operator in OCCS -

We know that,

$$\text{div. } \vec{A} = \vec{\nabla} \cdot \vec{A} = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u_1} (A_1 h_2 h_3) + \frac{\partial}{\partial u_2} (A_2 h_1 h_3) + \frac{\partial}{\partial u_3} (A_3 h_1 h_2) \right]$$

Let $A = \nabla \phi$ where ϕ is scalar field.

$$\text{div. } \vec{A} = \vec{\nabla} \cdot \nabla \phi = \nabla^2 \phi$$



Curl of Vector -

The curl of vector function obtained by Stokes' theorem,

$$\int_S (\nabla \times \vec{A}) \cdot d\vec{s} = \oint \vec{A} \cdot d\vec{l}$$

Using mean value -

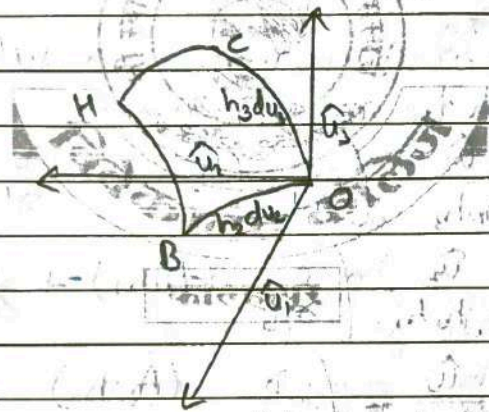
$$\text{curl } \vec{A} \int_S d\vec{s} = \oint \vec{A} \cdot d\vec{l} \text{ or } \text{curl } \vec{A} = \lim_{\delta s \rightarrow 0} \frac{\oint \vec{A} \cdot d\vec{s}}{\int_S d\vec{s}}$$

To evaluate this consider surface (s_1) OBHC

$$\oint_{OBHC} \vec{A} \cdot d\vec{l} = \oint_{OB} \vec{A} \cdot d\vec{l} + \oint_{BH} \vec{A} \cdot d\vec{l} + \oint_{HC} \vec{A} \cdot d\vec{l} + \oint_{CO} \vec{A} \cdot d\vec{l}$$

Let,

$$\vec{A} = A_1 \hat{u}_1 + A_2 \hat{u}_2 + A_3 \hat{u}_3$$



Evaluating integral,

$$\oint_{OBHC} \vec{A} \cdot d\vec{l} = A_2 h_2 du_2 + \left(A_3 h_3 du_3 + \frac{\partial A_3 h_3}{\partial u_2} du_2 du_3 \right) + \left(-A_2 h_2 du_2 - \frac{\partial A_2 h_2}{\partial u_3} du_2 du_3 \right) + (-A_3 h_3 du_3)$$

$$\oint_{OBHC} \vec{A} \cdot d\vec{l} = \left[\frac{\partial A_3 h_3}{\partial u_2} - \frac{\partial A_2 h_2}{\partial u_3} \right] du_2 du_3$$



Dividing by area of s_1 (OBHC) equal to $h_2 h_3 du_2 du_3$ and taking limit du_2 and du_3 approaches to zero, then component of $\text{curl } \vec{A}$ along \hat{u}_1 .

$$(\nabla \times \vec{A})_{u_1} = \frac{\left(\frac{\partial A_3 h_3}{\partial u_2} - \frac{\partial A_2 h_2}{\partial u_3} \right) du_2 du_3}{h_2 h_3 du_2 du_3}$$

$$= \frac{1}{h_1 h_3} \left(\frac{\partial A_3 h_3}{\partial u_2} - \frac{\partial A_2 h_2}{\partial u_3} \right)$$

Similarly by choosing areas S_2 and S_3 perpendicular to \hat{u}_2 and \hat{u}_3 at O respectively, we find component of curl \vec{A} along \hat{u}_2 and \hat{u}_3 . This leads to the result

$$(\nabla \times \vec{A})_{u_2} = \frac{1}{h_3 h_1} \left[\frac{\partial (A_1 h_1)}{\partial u_3} - \frac{\partial (A_3 h_3)}{\partial u_1} \right]$$

$$(\nabla \times \vec{A})_{u_3} = \frac{1}{h_1 h_2} \left[\frac{\partial (A_2 h_2)}{\partial u_1} - \frac{\partial (A_1 h_1)}{\partial u_2} \right]$$

From these three components, resultant equation for curl of vector field can be written as

$$\text{curl } \vec{A} = \nabla \times \vec{A} = \frac{\hat{u}_1}{h_2 h_3} \left[\frac{\partial (A_3 h_3)}{\partial u_2} - \frac{\partial (A_2 h_2)}{\partial u_3} \right]$$

$$+ \frac{\hat{u}_2}{h_3 h_1} \left[\frac{\partial (A_1 h_1)}{\partial u_3} - \frac{\partial (A_3 h_3)}{\partial u_1} \right]$$

$$+ \frac{\hat{u}_3}{h_1 h_2} \left[\frac{\partial (A_2 h_2)}{\partial u_1} - \frac{\partial (A_1 h_1)}{\partial u_2} \right]$$

$$\text{curl } \vec{A} = \nabla \times \vec{A} = \frac{1}{h_1 h_2 h_3} \begin{vmatrix} \hat{u}_1 h_1 & \hat{u}_2 h_2 & \hat{u}_3 h_3 \\ \frac{\partial}{\partial u_1} & \frac{\partial}{\partial u_2} & \frac{\partial}{\partial u_3} \\ A_1 h_1 & A_2 h_2 & A_3 h_3 \end{vmatrix}$$

This is the required result.



Vivekananda College Kolhapur (Autonomous).
Department of Physics: Internal examination 2018-19
B.Sc. III Semester V
Subject: Classical mechanics

Marks: 20 (Each question carry one mark)

Time : 20 min

Q.1 Attempt any ONE

(10)

1. Obtain an expression for the curl of vector field in orthogonal curvilinear co-ordinates.
2. Obtain an expression for the divergence of vector field in orthogonal curvilinear co-ordinate system. Extend the above formula in spherical polar co-ordinate system.

Q.2 Attempt any TWO

(10)

1. Obtain Laplacian operator in orthogonal curvilinear co-ordinate. Extend the result in cylindrical co-ordinates.
2. Obtain an expression for gradient of a scalar field in orthogonal curvilinear co-ordinate system.
3. Describe spherical polar co-ordinate system.



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SUPPLIMENT

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20

Suppliment No. :

Roll No. : 8524

Class : B.Sc-III, Sem V

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Subject : Classical mechanics

Test / Tutorial No. : Internal Exam

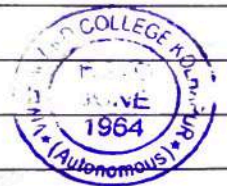
Div. :

Q2.

1. Equivalence of Lagrange and Newton's equations
⇒ In Newtonian mechanics the eqⁿ of motion involves vector quantities like force, momentum, which increase complexity in solving the problem. This approach can't avoid constraints present in problem.

These force of constraints, if not known make solution of problem difficult and even if they are known, use of rectangular / other commonly used co-ordinates may make solⁿ of problem impossible. These drawback removed in Lagrangian mechanics

where the technique involve scalar like potential and kinetic energies instead of vectors.



Q 2.

2. Write note on Degree of freedom and Constraint
⇒

Degrees of freedom -

The number of independent ways in which a mechanical system can move without violating any constraint which may be imposed on system is called number of degrees of freedom.

It is indicated by least possible number of co-ordinates to describe system completely

Eg - when single particle moves freely in space (x, y, z) it has three degrees of freedom

* Constraint -

when motion of particle / system of particle is restricted in some way then constraint have been introduced

Classification of constraint -

- (a) Scleronomous
- (b) Rheonomous
- (c) Holonomic
- (d) Non-holonomic



Q1.

2. Hamilton's Canonical equation of motion from Variational Principle

The Hamiltonian of system specified its total energy i.e kinetic energy and potential energy in terms of

$$\therefore L = T - V$$

$$\therefore H = \sum P_i \dot{q}_i - L$$

$$H = m \dot{q}_i q_i - (T - V)$$

$$= 2 \cdot \left(\frac{1}{2} m \dot{q}_i^2 \right) - (T - V)$$

$$= 2T - T + V$$

$$\therefore H = T + V$$

Hamilton Canonical Equation

$$\therefore H = H(q_i, P_i, t)$$

$$\therefore dH = \sum \frac{\partial H}{\partial q_i} dq_i + \sum \frac{\partial H}{\partial P_i} dP_i + \frac{\partial H}{\partial t} dt$$

5



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Suppliment No. :

Roll No. : 8523

Class : B.Sc-III, Sem V

Subject : classical mechanics

Test / Tutorial No. : Internal Exam

Div. : 2018-19

Q1.

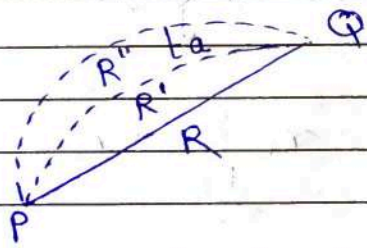
1. Hamilton Variational principle

⇒ Hamilton's principle -

The path actually transversed by Conservative holonomic dynamical system from time t_1 to t_2 is one over which the line integral of Lagrangian between limits t_1 and t_2 is stationary.

$$J = \int_{t_1}^{t_2} L dt = \text{extremum}$$

Let us consider that Conservative holonomic dynamical system moves from initial state P at t_1 to final state Q at t_2 .



To obtain Hamilton's principle following conditions must be satisfied

1) $\delta t = 0$ at the end points, i.e. at t_1 the particle must be at P and at t_2 the particle must be at Q

2) $\delta r = 0$ at the end points P and Q are fixed in space,

Acc. to Newton's 2nd law of motion the force F_i acting on i th particle of system

$$\therefore F_i = m_i \ddot{r}_i$$

\therefore From D'Alembert Principle,

$$\therefore \sum_{i=1}^n (F_i - m_i \ddot{r}_i) \delta r_i = 0$$

$$\therefore \sum_{i=1}^n F_i \delta r_i - \sum_{i=1}^n m_i \ddot{r}_i \delta r_i = 0 \quad \text{--- (2)}$$

But we have,

$$\ddot{r}_i \delta r_i = \frac{d}{dt} (\dot{r}_i \delta r_i) - \dot{r}_i \frac{d}{dt} (\delta r_i)$$

$$\therefore \frac{d}{dt} (\dot{r}_i \delta r_i) = \dot{r}_i \delta \dot{r}_i + \dot{r}_i \frac{d}{dt} (\delta r_i)$$

$$= \dot{r}_i \delta \dot{r}_i + \dot{r}_i \frac{d}{dt} (\delta r_i)$$



$$= \dot{r}_i \delta r_i + \delta \frac{1}{2} (v_i^2)$$

$$\therefore \dot{r}_i \delta r_i = \frac{d}{dt} (r_i \delta r_i) - \delta \frac{1}{2} (v_i^2)$$

The eqⁿ ② becomes,

$$\therefore \sum_{i=1}^n F_i \delta r_i - \left\{ \sum_{i=1}^n m_i \left[\frac{d}{dt} (r_i \delta r_i) - \delta \frac{1}{2} (v_i^2) \right] \right\} = 0$$

$$\therefore \sum_{i=1}^n F_i \delta r_i + \delta \sum_{i=1}^n \frac{1}{2} m_i v_i^2 = \sum_{i=1}^n \frac{d}{dt} (m_i v_i) \delta r_i$$

$$\therefore \sum_{i=1}^n F_i \delta r_i = \delta W = \text{work done by force } F_i \text{ during displacement} \quad \text{--- ③}$$

$$\therefore \sum_{i=1}^n \frac{1}{2} m_i v_i^2 = T = \text{K.E. of system}$$

\therefore Therefore eqⁿ ③ becomes,

$$\therefore \delta W + \delta T = \sum_i \frac{d}{dt} (m_i v_i) \delta r_i$$

\therefore Integrating above eqⁿ betⁿ limits t_1 and t_2

$$\begin{aligned} \therefore \int_{t_1}^{t_2} (\delta W + \delta T) dt &= \int_{t_1}^{t_2} \sum_{i=1}^n \frac{d}{dt} (m_i v_i) \delta r_i dt \\ &= \int_{t_1}^{t_2} \sum_{i=1}^n d(m_i v_i) \delta r_i \end{aligned}$$



$$= \left[\sum_{i=1}^n (m_i \dot{r}_i) \delta r_i \right]_P^Q$$

$$= 0$$

Since $\delta r_i = 0$ at the end points P and Q

$$\therefore \int_{t_1}^{t_2} (\delta W + \delta T) dt = 0 \quad \text{--- (4)}$$

But for conservative system, we have $\delta W = -\delta V$
 δV is change in potential energy.

$$\therefore \int_{t_1}^{t_2} (-\delta V + \delta T) dt = 0$$

$$\therefore \delta \int_{t_1}^{t_2} (T - V) dt = 0$$

$$\therefore \delta J = \delta \int_{t_1}^{t_2} L dt = 0$$

$$\therefore J = \int_{t_1}^{t_2} L dt = \text{Extremum}$$

10



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Suppliment No. :

Roll No. : 8510

Class : BSc-III Sem V

Subject : classical mechanics

Test / Tutorial No. : Internal Exam

Div. :

Q1

1. In newtonian mechanics the eqⁿ of motion involves vector quantities like force, momentum which increase complexity in solving the problem. This approach can't avoid constraints present in problems.

These force of constraints if not known make solution of problem difficult and even if they are known, use of rectangular and other commonly used co-ordinates may make solⁿ of problem impossible. These drawbacks removed in Lagrangian mechanics

where the technique involves scalars like potential and kinetic energies instead of vectors

5



2. ~~relation between H and L~~

2. Note on Degree of freedom and constraint

⇒ when motion of particle / system of particle is restricted in some way then constraints have been introduced

classification of constraints -

a. Scleronomic -

IF constraint relation do not explicitly depend time e.g. in case of rigid body.

b. rheonomic - IF constraint relation depend explicitly on time a bead sliding on moving wire in force free space.

c. Holonomic - Let r_1, r_2, \dots, r_n be the position co-ordinates of system and t denotes time then condition of all constraints are expressed as equations having the form $(r_1, r_2, \dots, r_n, t) = 0$

d. Non-holonomic - IF the condition of constraints are not so expressed as non-holonomic called as non-holonomic constraints.

5



Vivekananda College Kolhapur (Autonomous).
Department of Physics: Internal examination 2018-19

B.Sc. III Semester V

Subject: Quantum mechanics

Marks: 20 (Each question carry one mark)

Time : 20 min

Q.1. Long Answer question (Attempt any ONE) . [10]

- i) Obtain Schrodinger , s time independent equation and time dependent equation
- ii) Explain quantum mechanical treatment of linear harmonic oscillator and show that zero point energy of oscillator is $E_0 = \frac{1}{2} \hbar \omega$

Q.2. Short Answer question (Attempt any TWO). [10]

- i) Show that $[x, P_x] = i \hbar$ give its physical significance
- ii) Give physical significance of wave function
- iii) Obtain Schrodingers equation in spherical polar coordinate system for hydrogen atom



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Suppliment No. :

Roll No. : 8515

Class : B.Sc III

$\frac{-10}{20}$

Subject : Quantum mechanics

Test / Tutorial No. : Internal exam

Div. :

Q.1

i) a) Time-independent Schrodinger's equation:-

1) Total energy is given by

$$E = K.E + P.E.$$

$$E = \frac{1}{2}mv^2 + V$$

$$E = \frac{1}{2} \frac{m^2v^2}{m} + V$$

$$E = \frac{p^2}{2m} + V$$

$$p^2 = 2m(E - V)$$

$$p = [2m(E - V)]^{1/2} \quad \text{--- (1)}$$

$$\therefore \text{de-Broglie wavelength } \lambda = \frac{h}{p} = \frac{h}{[2m(E - V)]^{1/2}}$$

e) He assumed that the wave functⁿ ψ is governed by the equation

$$\psi = \psi_0 e^{-i\omega t} \quad \text{--- (2)}$$

Wave equation in Cartesian co-ordinates is given by



$$\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} = \frac{1}{u^2} \frac{\partial^2 \psi}{\partial t^2} \quad \text{--- (3)}$$

Differentiating eqⁿ (3) twice with respect to t we get

$$\frac{\partial^2 \psi}{\partial t^2} = -\omega^2 \psi_0 e^{-i\omega t} = -\omega^2 \psi$$

Substituting this value in eqⁿ 3

$$\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} = -\left(\frac{\omega^2}{u^2}\right) \psi$$

$$\frac{\omega^2}{u^2} = \frac{(2\pi\nu)^2}{(2\lambda)^2} = \frac{4\pi^2}{\lambda^2}$$

$$\therefore \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} = -\left(\frac{4\pi^2}{\lambda^2}\right) \psi$$

$$\nabla^2 \psi + \frac{4\pi^2}{\lambda^2} \psi = 0 \quad \text{--- (4)}$$

Eqⁿ (4) is general eqⁿ and is independent of time

Substituting this value of λ in eqⁿ 4

$$\nabla^2 \psi + \frac{4\pi^2 [2m(E-V)]}{h^2} \psi = 0$$

$$\nabla^2 \psi + \left(\frac{8\pi^2 m}{h^2}\right) (E-V) \psi = 0$$

Also, $\hbar = \frac{h}{2\pi}$

$$\therefore \nabla^2 \psi + \left(\frac{2m}{\hbar^2}\right) (E-V) \psi = 0 \quad \text{--- (5)}$$

This is the Schrodinger's time-independent wave eqⁿ



b) Time-dependent Schrodinger's eqⁿ

The wave functⁿ

$$\psi = \psi_0 e^{-i\omega t}$$

$$\frac{\partial \psi}{\partial t} = -i\omega \psi_0 e^{-i\omega t}$$

$$\frac{\partial \psi}{\partial t} = -i(2\pi\nu)\psi$$

$$\frac{\partial \psi}{\partial t} = -i(2\pi) \left(\frac{E}{h}\right) \psi$$

$$E\psi = -i \left(\frac{h}{2\pi}\right) \frac{\partial \psi}{\partial t}$$

$$E\psi = i\hbar \frac{\partial \psi}{\partial t} \quad \left(\hbar = \frac{h}{2\pi} \text{ and } i^2 = -1 \right)$$

(we know)

$$\nabla^2 \psi + \frac{2m}{\hbar^2} E\psi - \frac{2m}{\hbar^2} V\psi = 0$$

Substituting value of $E\psi = i\hbar \frac{\partial \psi}{\partial t}$ in above eqⁿ

$$\nabla^2 \psi + \frac{2m}{\hbar^2} i\hbar \frac{\partial \psi}{\partial t} - \frac{2m}{\hbar^2} V\psi = 0$$

$$\nabla^2 \psi - \left(\frac{2m}{\hbar^2}\right) V\psi = -\left(\frac{2m}{\hbar^2}\right) i\hbar \frac{\partial \psi}{\partial t}$$

or

$$-\left(\frac{\hbar^2}{2m}\right) \nabla^2 \psi + V\psi = i\hbar \frac{\partial \psi}{\partial t}$$

$$\left(\frac{-\hbar^2}{2m} \nabla^2 + V\right) \psi = i\hbar \frac{\partial \psi}{\partial t}$$

$$H\psi = E\psi$$

This is the Schrodinger's time-dependent wave eqⁿ



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Subject : Quantum Mechanics.

Test / Tutorial No. : Internal Exam

Div. :

Suppliment No. :

Roll No. : 8519

Class : B.Sc. III, Sem-V (18-19)

13
20

Q.2)

1) Consider the action of commutator $[x, P_x]$ on wave $\psi(x, y, z)$.

$$[x, P_x] \psi = [x, -i\hbar \frac{\partial}{\partial x}] \psi$$

$$= -i\hbar [x \frac{\partial}{\partial x}] \psi$$

$$= -i\hbar \left(x \frac{\partial}{\partial x} - \frac{\partial}{\partial x} x \right) \psi$$

$$= -i\hbar \left\{ x \frac{\partial \psi}{\partial x} - \frac{\partial}{\partial x} (x\psi) \right\}$$

$$= -i\hbar \left\{ x \frac{\partial \psi}{\partial x} - x \frac{\partial \psi}{\partial x} - \psi \right\}$$

$$= -i\hbar (-\psi)$$

$$= i\hbar \psi$$

$$[x, P_x] \psi = i\hbar \psi$$

$$\therefore [x, P_x] = i\hbar$$



$$\nabla^2 \psi + \frac{4\pi^2}{\lambda^2} \psi = 0 \quad (5)$$

Eqⁿ (5) is general eqⁿ & is independent of time.

By the concept of wave mechanics de-Broglie wavelength

$$\lambda = \frac{h}{p} = \frac{h}{[2m(E-V)]^{1/2}} \quad (\text{from eqⁿ (1)})$$

Putting this value of λ in eqⁿ (5) we get,

$$\nabla^2 \psi + \frac{4\pi^2 [2m(E-V)]}{h^2} \psi = 0$$

$$\nabla^2 \psi + \left(\frac{8\pi^2 m}{h^2} \right) (E-V) \psi = 0 \quad (6)$$

Also, $h = \frac{h}{2\pi}$

$$\therefore \nabla^2 \psi + \left(\frac{2m}{\hbar^2} \right) (E-V) \psi = 0 \quad (7)$$

Eqⁿ (7) represents Schrödinger's time independent wave eqⁿ.



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Suppliment No. :

Roll No. : 8505

Class : B.Sc III, Sem-V (18-19)

Subject : Quantum Mechanics.

Test / Tutorial No. : Internal Exam

Div. :

20
20

Q.1)

1) (a) Time independent Schrödinger's Equation-

Assumed that de-Borglie wavelength for any particle moving in a field of force with potential energy V . Therefore total Energy is given by.

$$E = KE + PE$$

$$= \frac{1}{2} m v^2 + V$$

$$= \frac{1}{2} \frac{m^2 v^2}{m} + V$$

$$E = \frac{p^2}{2m} + V$$

$$p^2 = 2m(E - V)$$

$$p = [2m(E - V)]^{1/2} \quad (1)$$

$$\therefore \text{De-Borglie wavelength } \lambda = \frac{h}{p} = \frac{h}{[2m(E - V)]^{1/2}} \quad (2)$$

The wave function ψ is governed by the equation

$$\psi = \psi_0 e^{-iwt} \quad (3)$$

The wave equation in cartesian co-ordinates is given by,

$$\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} = \frac{1}{\psi} \frac{\partial^2 \psi}{\partial t^2} \quad (4)$$

Now, differentiate eqn (3) twice w.r. to t we get,



$$\frac{\partial^2 \psi}{\partial t^2} = -\omega^2 \psi_0 e^{-i\omega t} = -\omega^2 \psi$$

Putting this value in eqⁿ (4) we get,

$$\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} = -\left(\frac{\omega^2}{v^2}\right) \psi$$

where, $\frac{\omega^2}{v^2} = \frac{(2\pi\nu)^2}{(v\lambda)^2} = \frac{4\pi^2}{\lambda^2}$

$$\therefore \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} = -\left(\frac{4\pi^2}{\lambda^2}\right) \psi$$

$$\therefore \nabla^2 \psi + \frac{4\pi^2}{\lambda^2} \psi = 0 \quad \text{--- (5)}$$

Eqⁿ (5) is general eqⁿ & is independent of time.

By the concept of wave mechanics, de-Broglie wavelength

$$\lambda = \frac{h}{p} = \frac{h}{[2m(E-V)]^{1/2}} \quad (\text{from eqⁿ (1)})$$

Putting this value of λ in eqⁿ (5) we get,

$$\nabla^2 \psi + \frac{4\pi^2 [2m(E-V)] \psi}{h^2} = 0$$

$$\nabla^2 \psi + \left(\frac{8\pi^2 m}{h^2}\right) (E-V) \psi = 0 \quad \text{--- (6)}$$

Also, $\hbar = \frac{h}{2\pi}$

$$\therefore \nabla^2 \psi + \left(\frac{2m}{\hbar^2}\right) (E-V) \psi = 0 \quad \text{--- (7)}$$

Eqⁿ (7) represents Schrödinger's time independent wave eqⁿ.

(b) Time dependent Schrödinger's equation →

The wave function, $\psi = \psi_0 e^{-i\omega t}$

$$\frac{\partial \psi}{\partial t} = -i\omega \psi_0 e^{-i\omega t}$$

$$\frac{\partial \psi}{\partial t} = -i(2\pi\nu) \psi$$



$$\frac{\partial \psi}{\partial t} = -i(2\pi) \left(\frac{E}{h} \right) \psi$$

$$E\psi = -i \left(\frac{h}{2\pi} \right) \frac{\partial \psi}{\partial t}$$

$$E\psi = i\hbar \frac{\partial \psi}{\partial t} \quad \left(\because \hbar = \frac{h}{2\pi} \text{ \& } i^2 = -1 \right) \quad (8)$$

From eqⁿ (7)

$$\nabla^2 \psi + \frac{2m}{\hbar^2} E\psi - \frac{2m}{\hbar^2} V\psi = 0 \quad (9)$$

Putting the value of $E\psi = i\hbar \frac{\partial \psi}{\partial t}$ in eqⁿ (9) we get,

$$\nabla^2 \psi + \frac{2m}{\hbar^2} \frac{\partial \psi}{\partial t} - \frac{2m}{\hbar^2} V\psi = 0$$

$$\nabla^2 \psi - \left(\frac{2m}{\hbar^2} \right) V\psi = \left(\frac{2m}{\hbar^2} \right) i\hbar \frac{\partial \psi}{\partial t}$$

$$\left(-\frac{\hbar^2}{2m} \right) \nabla^2 \psi + V\psi = i\hbar \frac{\partial \psi}{\partial t} \quad (10)$$

Eqⁿ (10) Represents Schrödinger's time dependent wave eqⁿ.

Q.2)

1) Consider the action of commutator $[x, P_x]$ on wave fn $\psi(x, y, z)$

$$[x, P_x] \psi = [x, -i\hbar \frac{\partial}{\partial x}] \psi$$

$$= -i\hbar \left[x, \frac{\partial}{\partial x} \right] \psi$$

$$= -i\hbar \left(x \frac{\partial}{\partial x} - \frac{\partial}{\partial x} x \right) \psi$$

$$[x, P_x] \psi = -i\hbar \left\{ x \frac{\partial \psi}{\partial x} - \frac{\partial}{\partial x} (x\psi) \right\}$$

$$= -i\hbar \left\{ x \frac{\partial \psi}{\partial x} - x \frac{\partial \psi}{\partial x} - \psi \right\}$$

$$= -i\hbar (-\psi)$$



$$\therefore [x, P_x] \Psi = i\hbar \Psi$$

$$\therefore [x, P_x] = i\hbar$$

Q.2)

3) The variable quantity or oscillatory function characterising, de-Broglie wave is called as wave function denoted by symbol Ψ . It is always associated with moving particle. The value of wave function associated with moving particle at a particular point x, y, z in space and the time t is related to the finding a particle at that time Ψ .

i. Ψ must be finite for all values of x, y and z .

ii. Ψ must be single valued i.e. for each set of values of x, y and z , Ψ must have only one value.

iii. Ψ must be continuous in all regions, except in those regions where the potential energy $V(x, y, z) = \infty$

iv. The partial derivatives of Ψ i.e. $\frac{\partial \Psi}{\partial x}$, $\frac{\partial \Psi}{\partial y}$, $\frac{\partial \Psi}{\partial z}$ must also

be finite, single-valued and continuous at all points, except at points where the potential $V(x)$ is infinite.

v. Ψ must vanish at infinity i.e. $\Psi = 0$, as $x \rightarrow \pm\infty$ or $z \rightarrow \pm\infty$.

